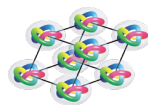


クォーク・グルーオンプラズマの物理とその周辺

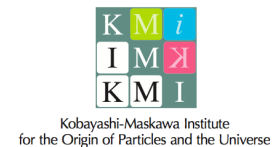
Department of physics, Hiroshima University
International Institute for Sustainability with Knotted Chiral Meta Matter / SKCM².

Hiroshima University
Kobayashi Maskawa Institute, Nagoya University

Chiho NONAKA



SKCM²
WPI HIROSHIMA UNIVERSITY



August 5, 2025@核融合研

クォーク・グルーオン プラズマの物理



- 世界規模の実験が稼働中：理論と実験の検証
- 周辺物理との関連
 - 理論的概念、ツール：共通点
 - 基礎理論
 - 相転移、相図
 - 平衡、非平衡物理
 - 熱力学量
 - ダイナミクスの記述
 - 相対論的流体模型
 - 分子動力学
 - Particle in Cell
 - データ解析、大規模計算
 - ベイジアン解析



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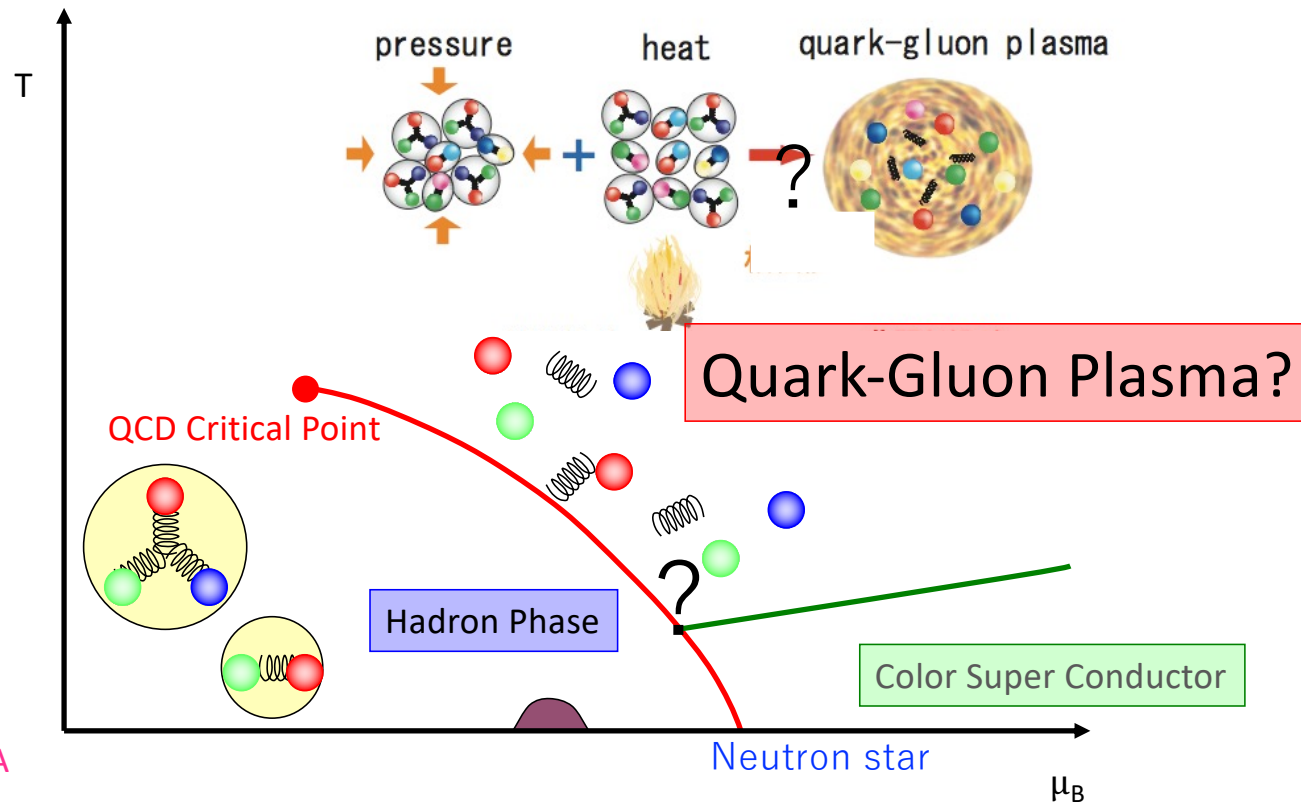
クォークとグルーオンの多体系の物理



• Quarks and gluons at extreme conditions

Quark-Gluon Plasma

- High temperature and/or high density

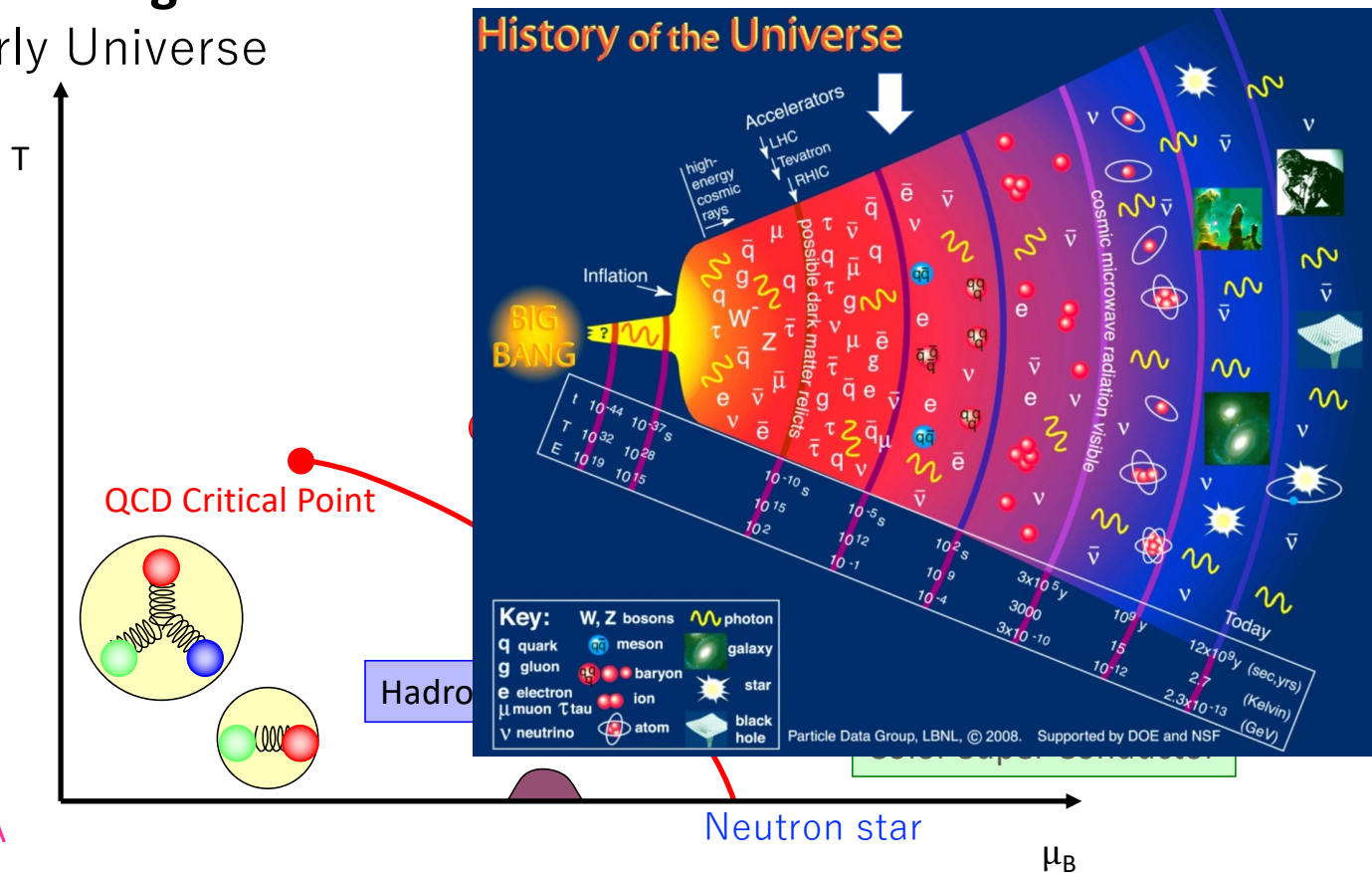


What is the QGP?



- Quarks and gluons at extreme conditions
 - Early Universe

Quark-Gluon Plasma



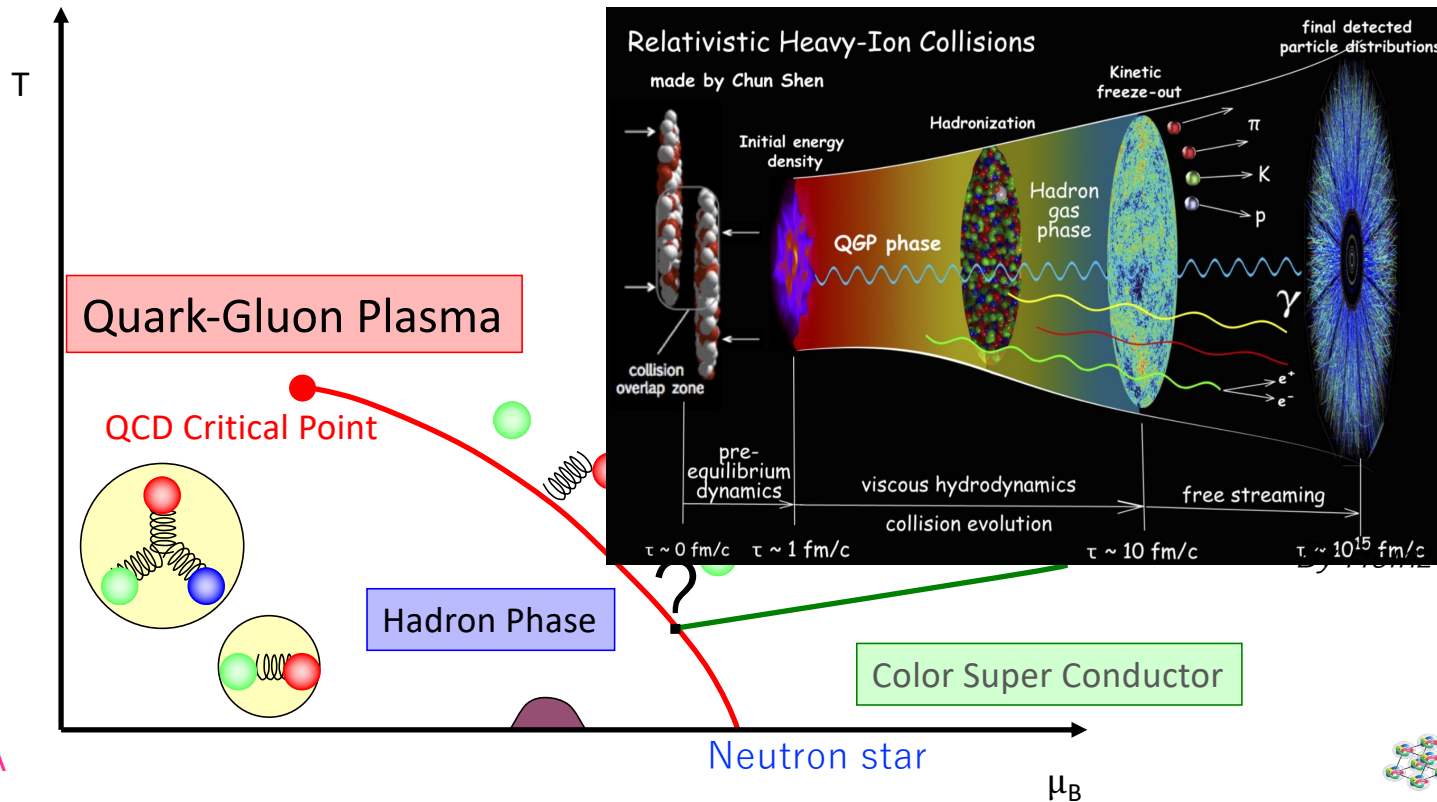
What is the QGP?



• Quarks and gluons at extreme conditions

Quark-Gluon Plasma

- Relativistic Heavy Ion Collisions : Little Bang



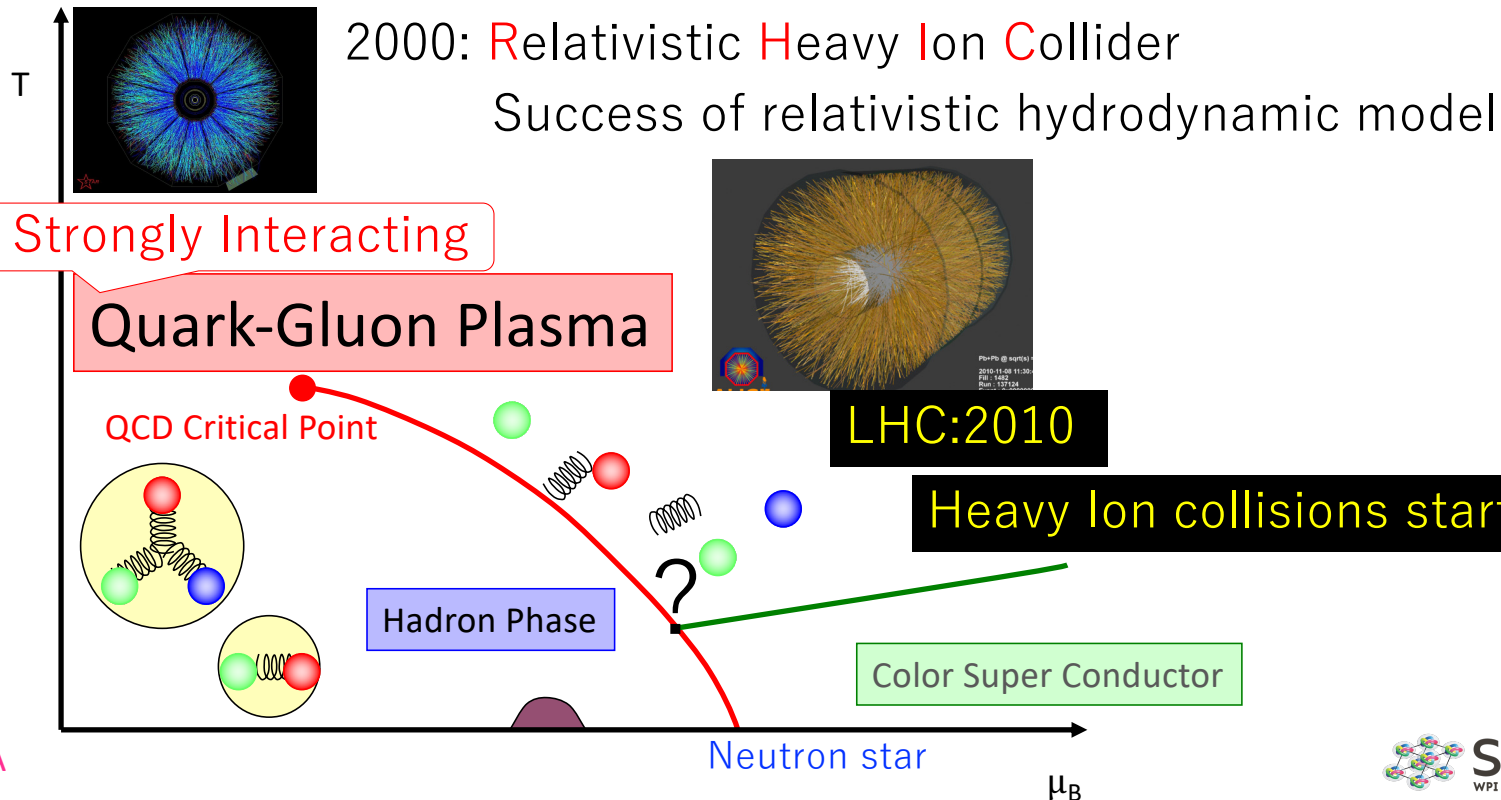
What is the sQGP?



- **Quarks and gluons at extreme conditions**

Quark-Gluon Plasma

- Relativistic Heavy Ion Collisions

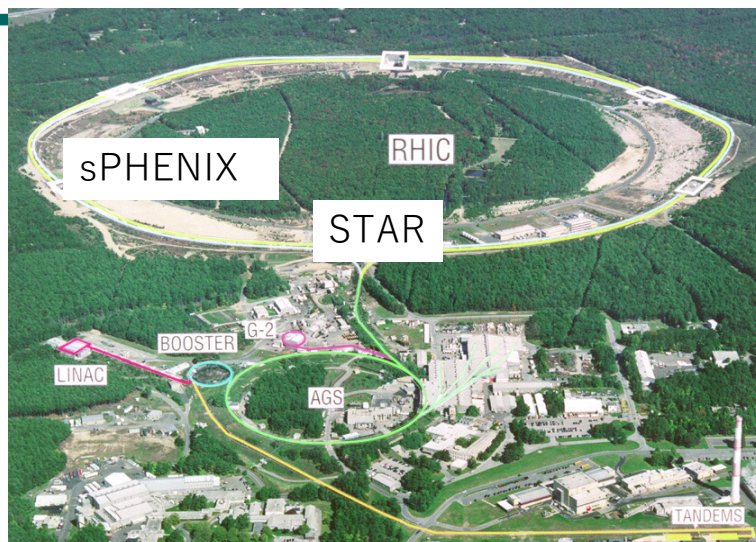


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Heavy Ion Collisions



RHIC@BNL

Large Hadron Collider@CERN

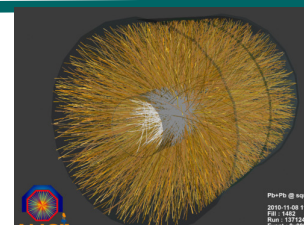
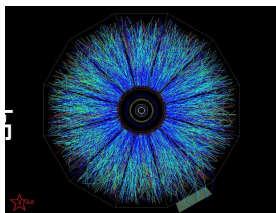


Heavy Ion Collisions



STAR@RHIC

ALICE@LHC



Au+Au(Beam Energy Scan)
7.7, 11.5, 19.8, 27, 39

p+p,
d+Au, He+Au
U+U, Au+Au,
200

p+Pb
Pb+Pb
2760 5020

GeV

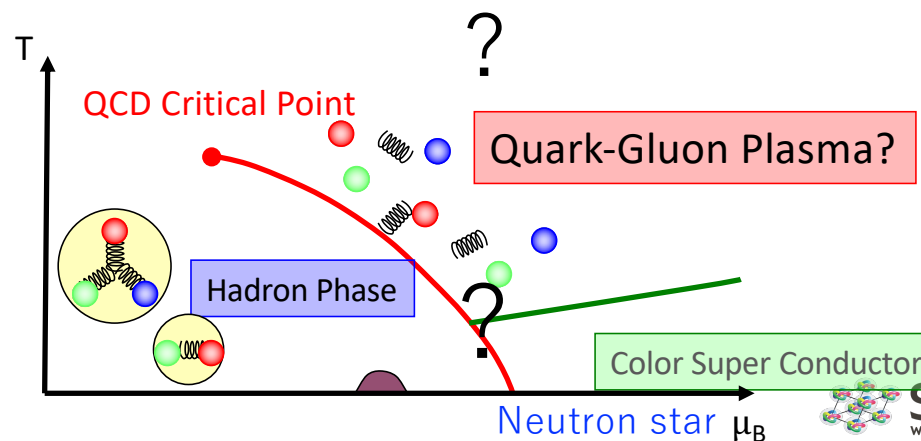
RHIC

LHC

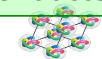
$\sqrt{s_{NN}}$

- 様々な衝突エネルギー
- 様々なものを衝突させている。

QCD 相図のどこに対応
するのか？



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統計模型

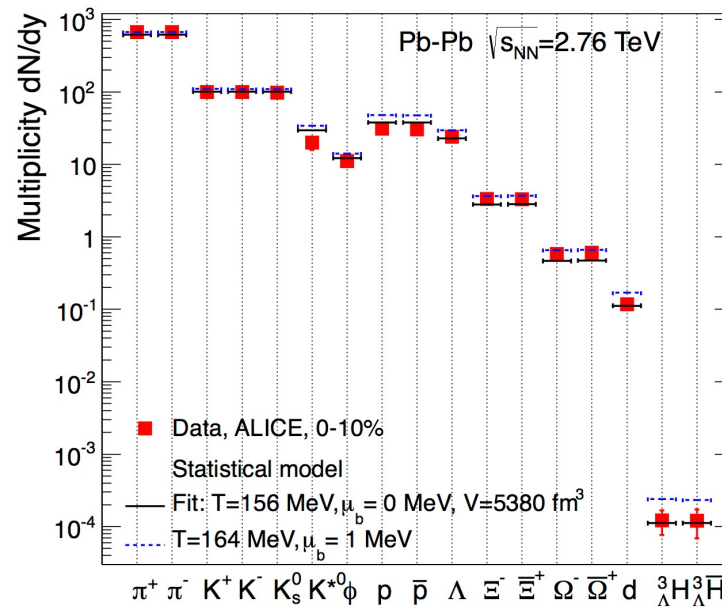


	p+p, d+Au, He+Au	p+p p+Pb	
Au+Au(Beam Energy Scan)	U+U, Au+Au,	Pb+Pb	Pb+Pb
7.7, 11.5, 19.8, 27, 39	200	2760	5020
			GeV

実験データ



QCD相図のどこを見ているのか? $\sqrt{s_{NN}}$



$$n_i = N_i/V = g_i \int \frac{d^3p}{(2\pi)^3} \frac{1}{e^{-(\epsilon_i(p)-\mu)/T} \pm 1}$$

Pb+Pb 2760 GeV
 $\rightarrow \sim T=156$ MeV, $\mu=0$



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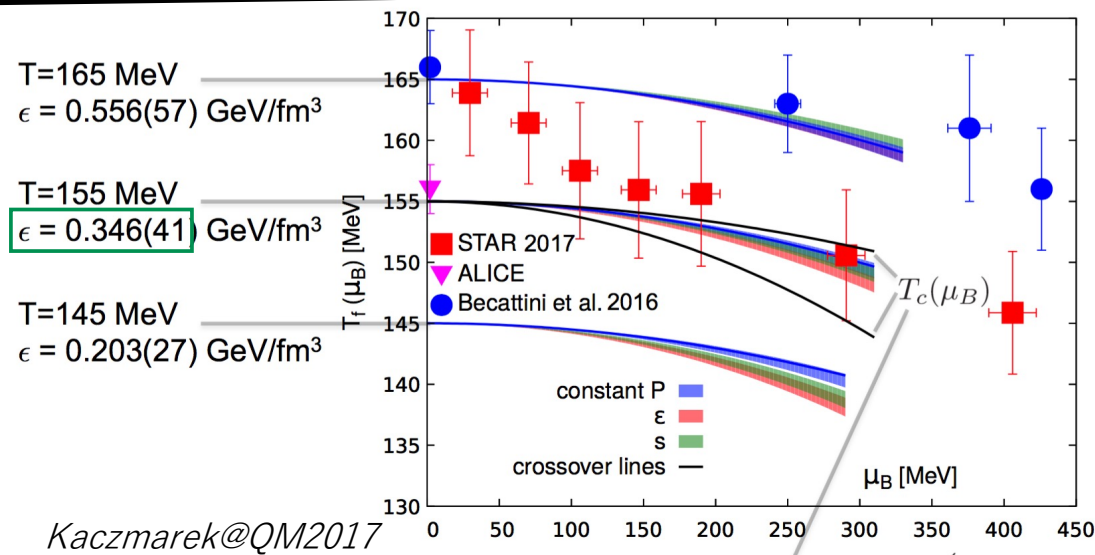
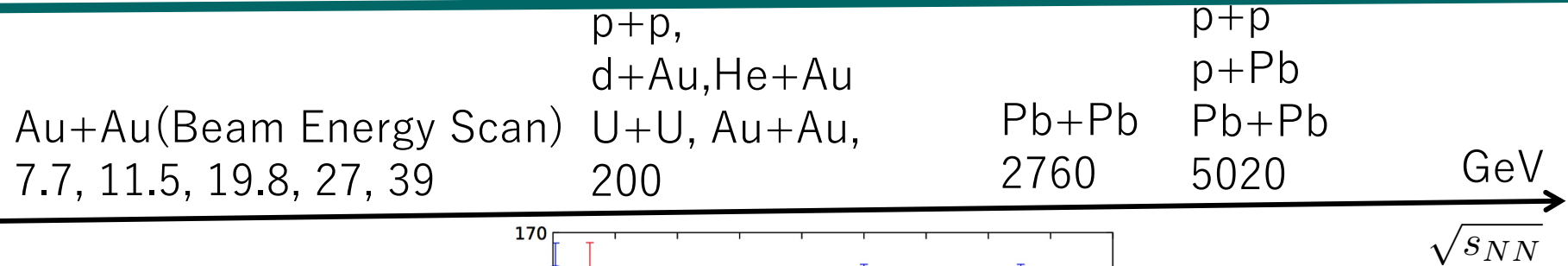
Ex: Stachel et al, arXiv:1311.4662



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実験とQCD相図



格子 QCD
(第一原理計算)
 $T_c = 154(9)$ MeV
HotQCD
PRD85,054503
(2012)

compare well with estimates of the crossover line: $T_c(\mu_B) = T_c(0) \left(1 - \kappa_2^c \left(\frac{\mu_B}{T_c(0)} \right)^2 \right)$



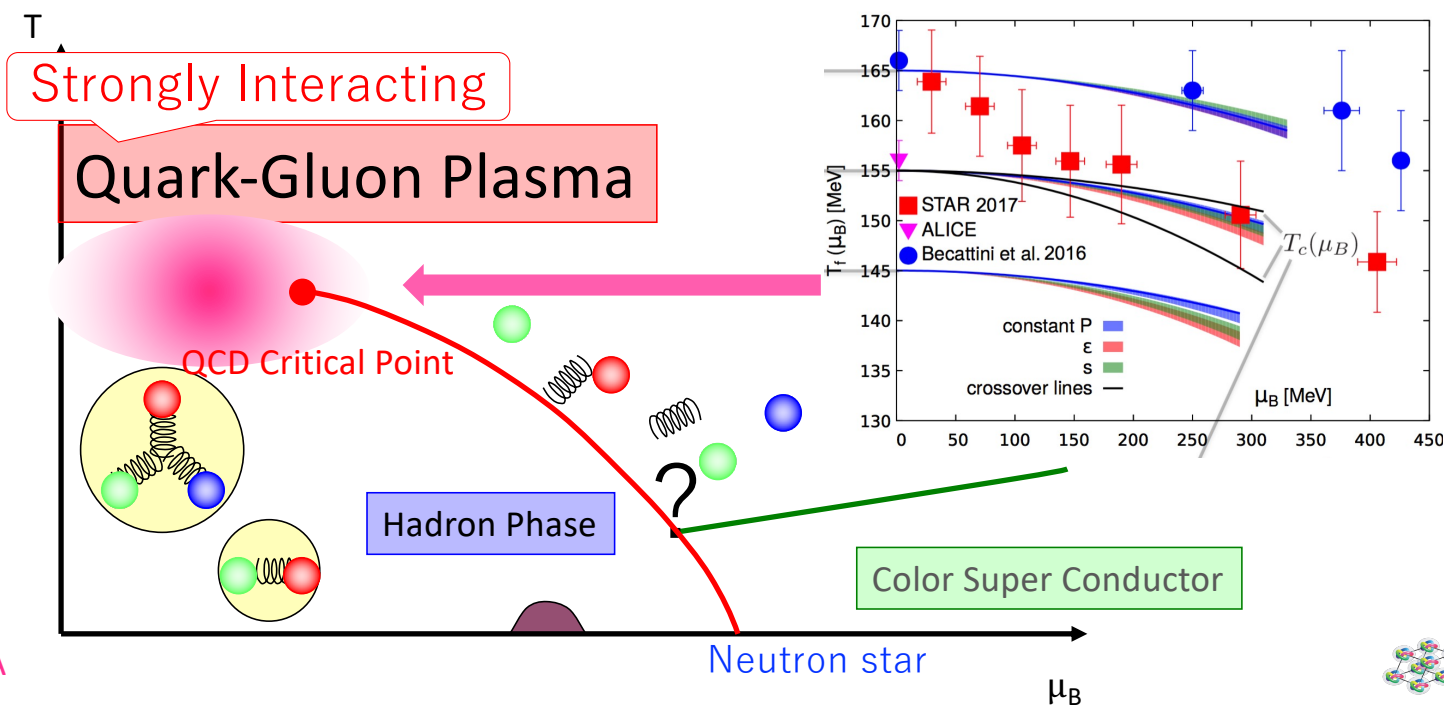
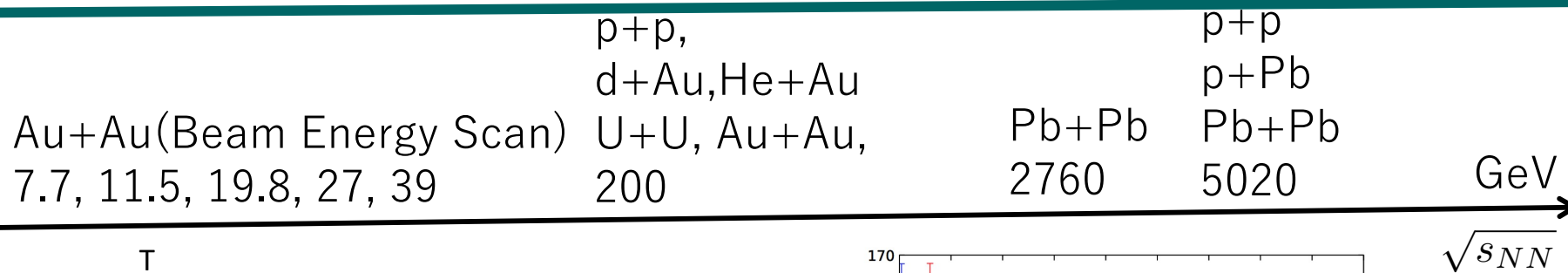
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実験とQCD相図



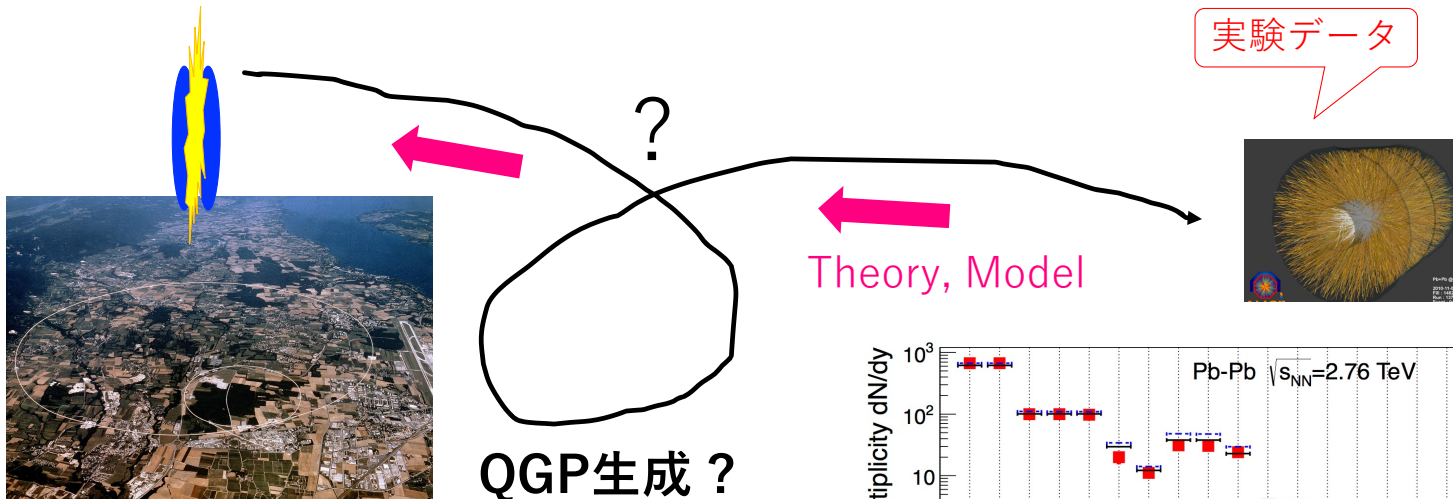
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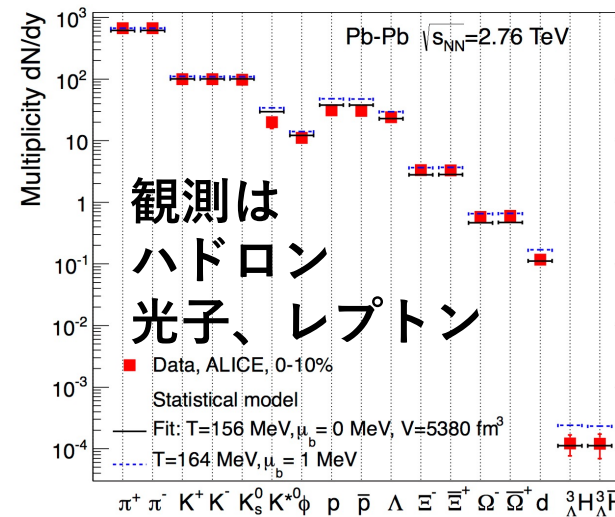


QGP生成？



模型構築

- QGP生成のメカニズム
- 時間発展
- QGP, ハドロンの性質



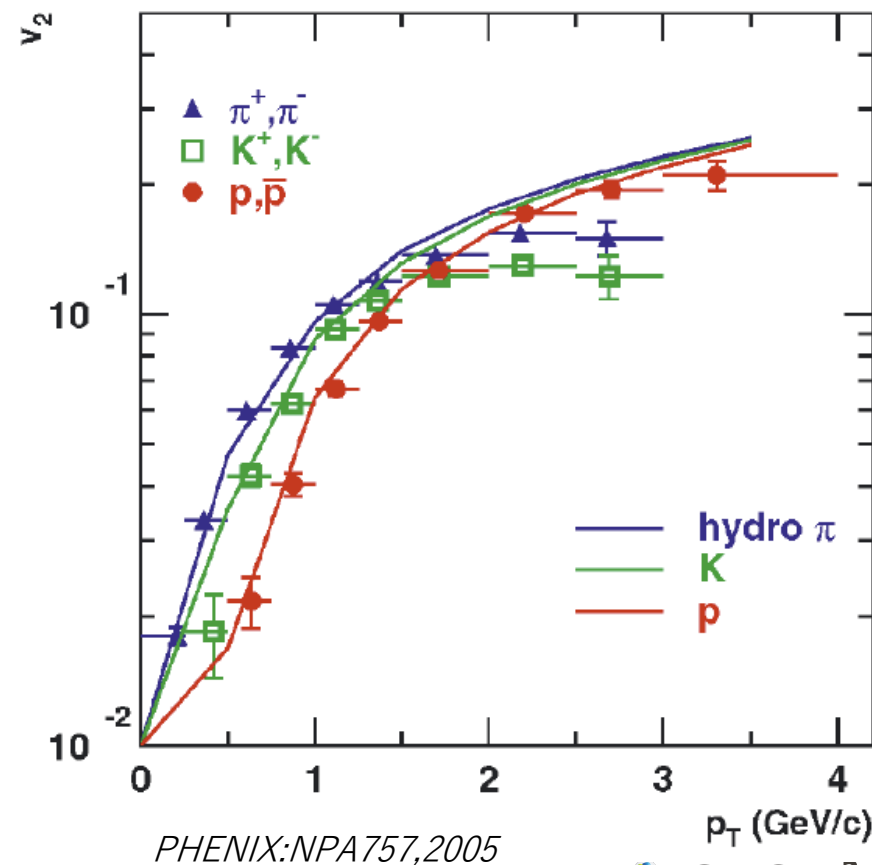
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相対論的流体模型



- 2001年RHIC稼働後標準的な模型に
強い楕円フロー
- 衝突後の粒子多重発生
流体描像で取り扱う
Landau
Bjorken

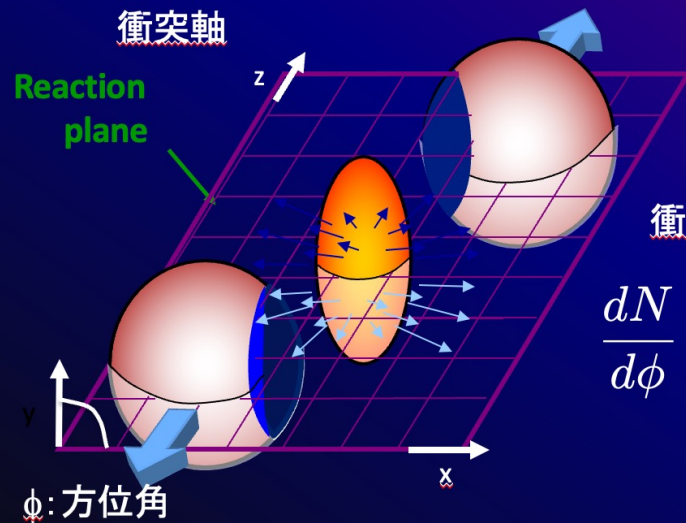
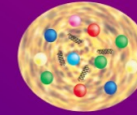


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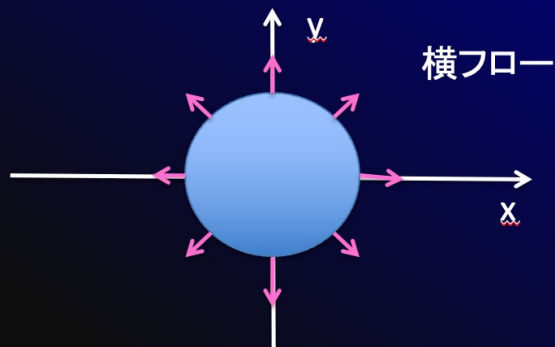
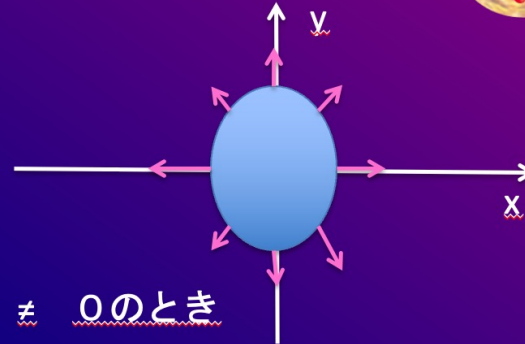
粒子の集団運動：楕円フロー



衝突係数 $\neq 0$ のとき

$$\frac{dN}{d\phi} \sim N_0(1 + 2v_1 \cos \phi + 2v_2 \cos 2\phi)$$

直接 楕円

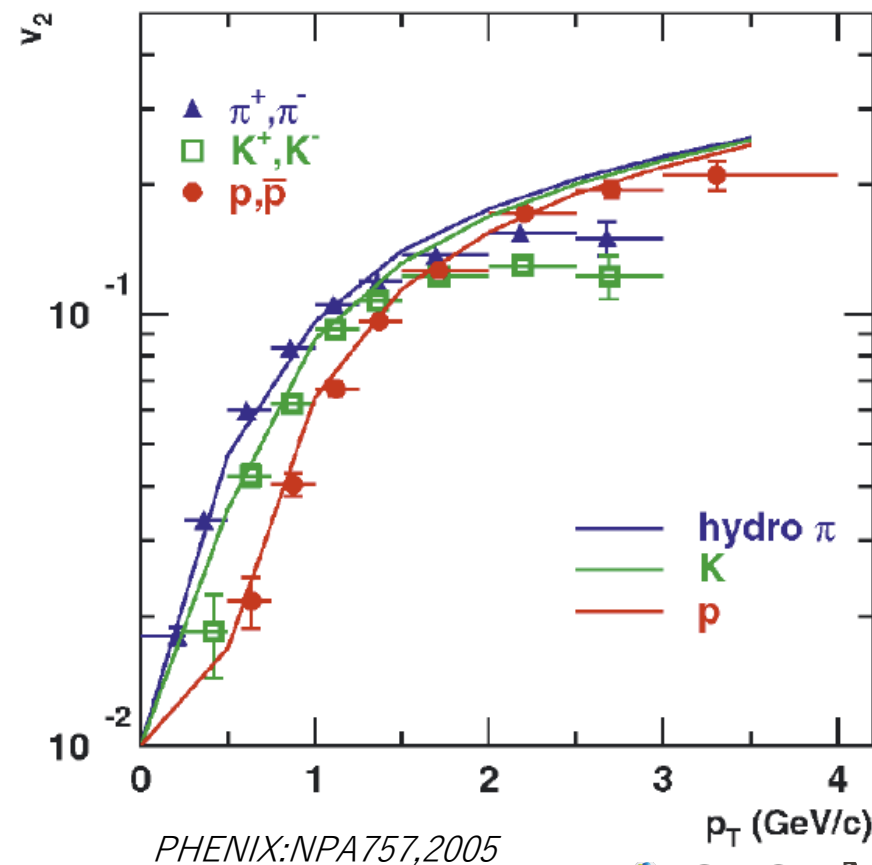


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相対論的流体模型



- 2001年RHIC稼働後標準的な模型に
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流体描像で取り扱う
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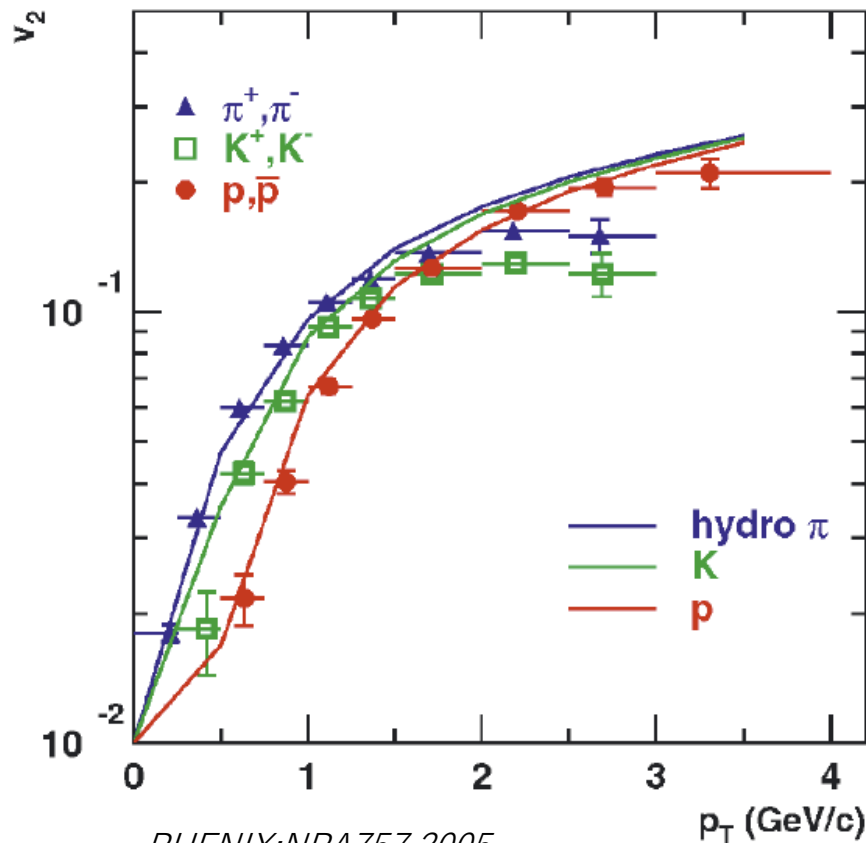
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相対論的流体模型



PHENIX:NPA757,2005

特徴

- 楕円フローの質量依存性
粒子の集団運動

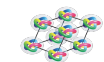
- 衝突直後の熱平衡化
流体の初期時刻
 $t_0 \sim 0.6 \text{ fm}$

$$1 \text{ fm} = 10^{-15} \text{ m}$$

問題点も存在→現在はほとんど解決



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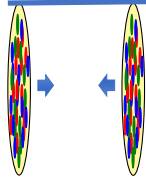


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Hydrodynamic Model

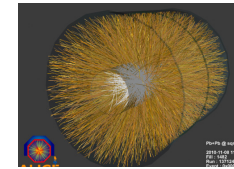


collisions



QGP
Production
?

Experimental data



Hadrons
Leptons
Photons

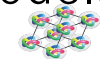
Multiple particle production

➡ Hydrodynamic picture
Landau
Bjorken

Success of hydrodynamic model at RHIC
Relativistic viscous hydrodynamic model
One of important phenomenological models



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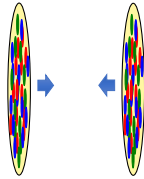
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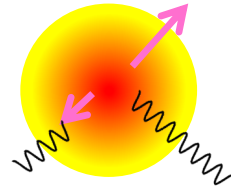
Equation of State



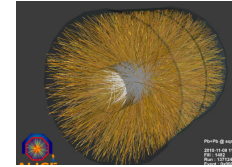
collisions



hydrodynamics



Experimental data



Hydrodynamics

QCD phase
diagram
EoS: lattice QCD

$$\partial_\mu T^{\mu\nu} = 0$$

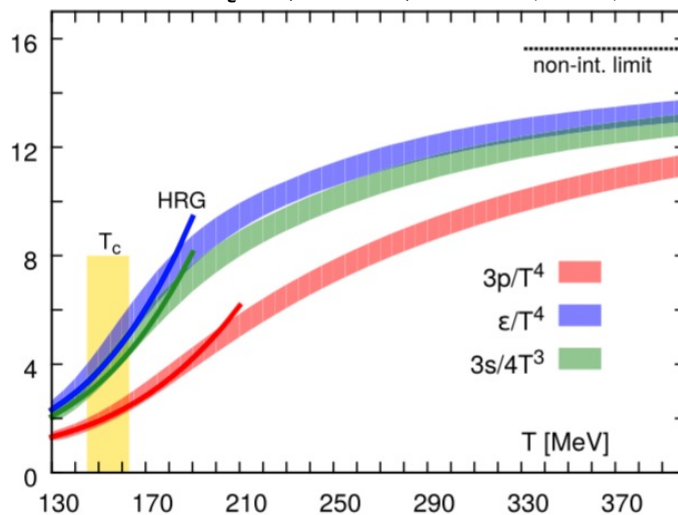
Equation of State



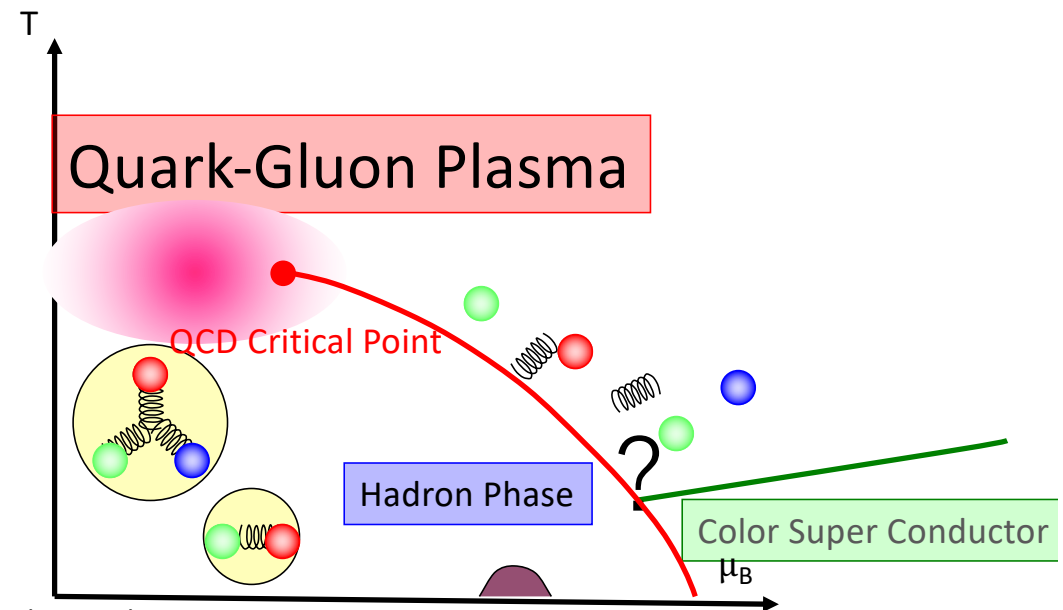
• Equation of State

• Lattice QCD

HotQCD, PRD90, 094503 (2014)



(2+1) flavor, Highly improved staggered quark action
 $N_t=6,8,10,12, N_s=4N_t \rightarrow$ continuum limit
 Parametrization of EoS



Neutron star



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$T_c \sim 155$ MeV

finite μ : sign problem



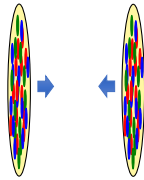
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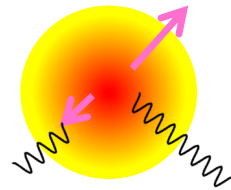
Equation of State



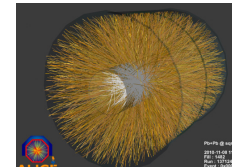
collisions



hydrodynamics



Experimental data



Hydrodynamics

QCD phase diagram

EoS: lattice QCD

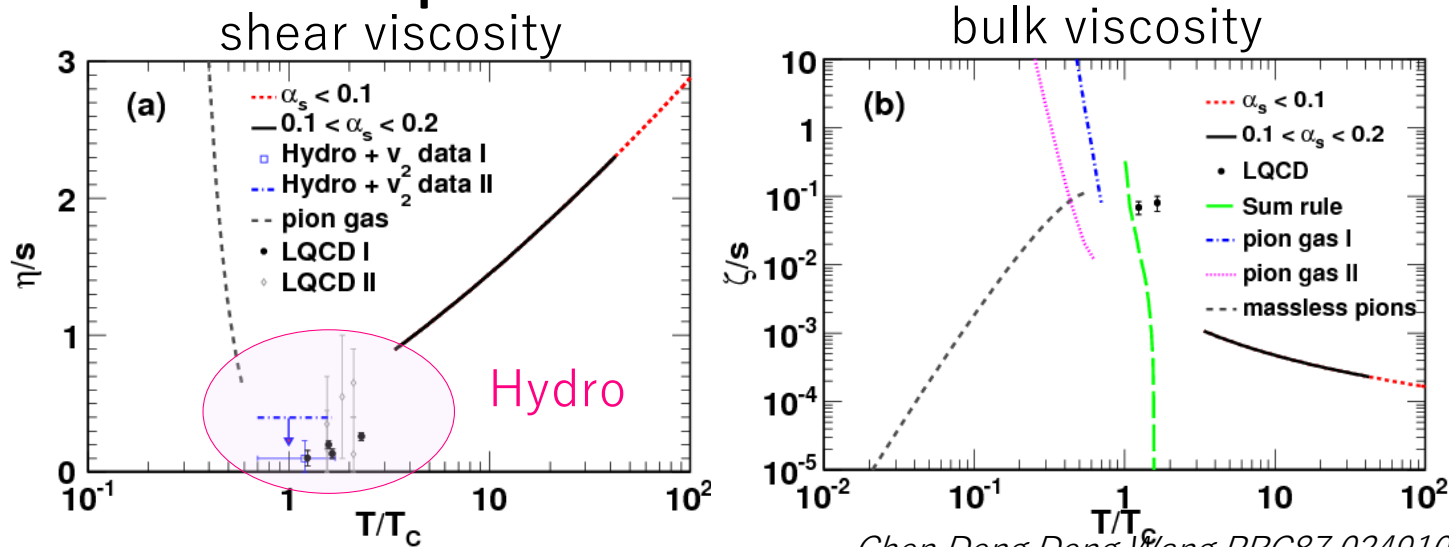
Shear and bulk viscosities

$$\partial_\mu T^{\mu\nu} = 0$$

Property of QGP



• Current Status for transport coefficients



- Shear viscosity takes the minimum around T_c . Cf. $\eta/s = 1/4\pi$ AdS/CFT
- Hydrodynamic model constant η/s

- Chen, Deng, Dong, Wang, PRC87, 024910 (2013)
- Bulk viscosity Temperature dependence is unclear.
 - Hydrodynamic model vanishing

パラメータを仮定 + ベイジアン解析



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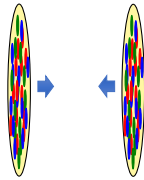
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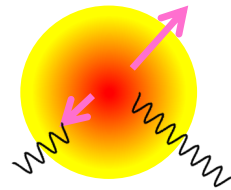
Equation of State



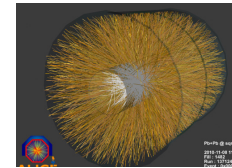
collisions



hydrodynamics



Experimental data



Hydrodynamics

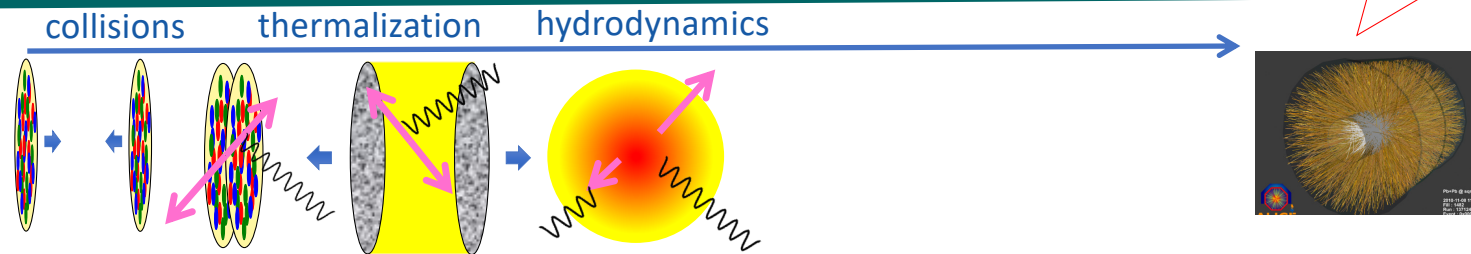
QCD phase diagram

EoS: lattice QCD

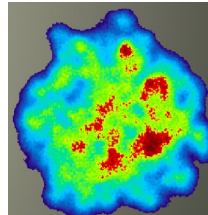
Shear and bulk viscosities

$$\partial_\mu T^{\mu\nu} = 0$$

Initial Condition



Initial
conditions



Energy (entropy)
density distributions

最も難しい問題の一つ

Hydrodynamics

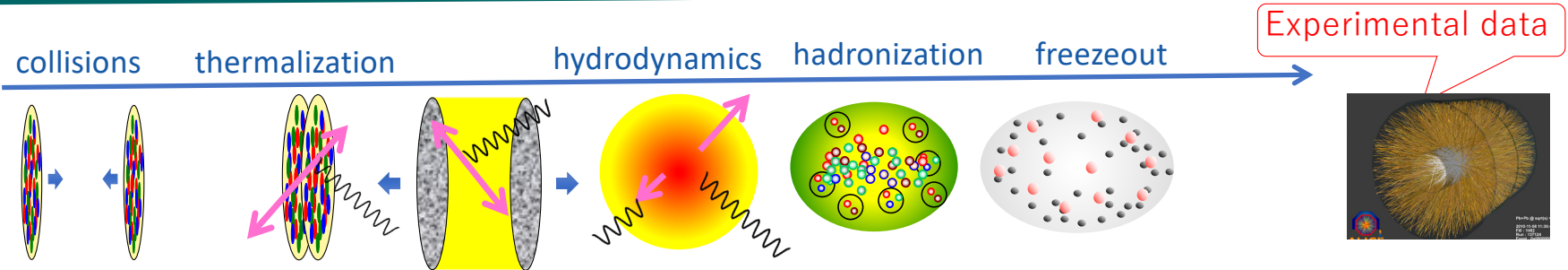
QCD phase diagram

EoS: lattice QCD

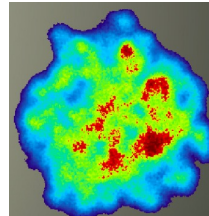
Shear and bulk viscosities

$$\partial_\mu T^{\mu\nu} = 0$$

From Fluid to Particle



Initial
conditions



Energy (entropy)
density distributions

Hydrodynamics

QCD phase
diagram
EoS: lattice QCD
Shear and bulk
viscosities

Hadronization

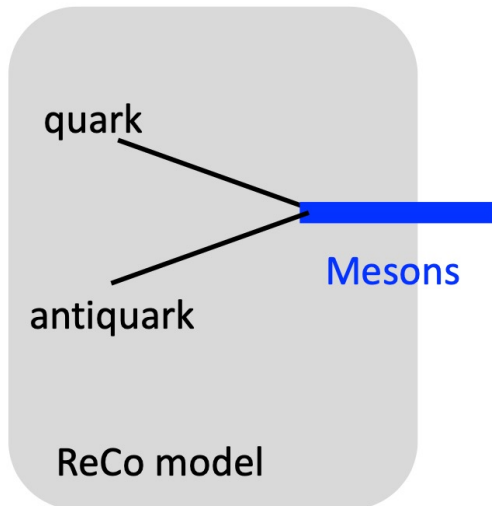
Recombination Model

Recombination Model

Fries, Mueller, CN and Bass, PRC68(2003)

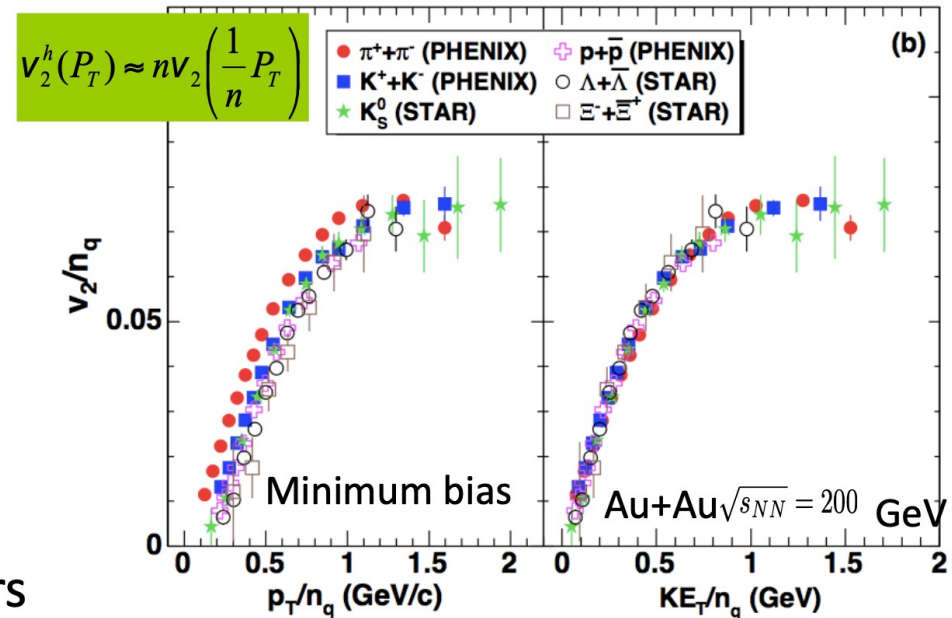


- One of successful models



- Baryon/Meson ratios
- Nuclear modification factors
- Quark number scaling in elliptic flow

Ex: Quark number scaring



PHENIX, nucl-ex:0608033



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Radiative Recombination in QGP

One of Photon Production Processes

Fujii, Itakura, CN, NPA 967 (2017)

Fujii, Itakura, Miyachi, CN,

Phys.Rev.C 106 (2022) 3, 034906

- Photon emission at hadronization process
 - Photon's flow is as strong as hadrons' flow.
- A photon is produced from pairing of quarks
 - Radiative recombination brings enhancement of photon yield.

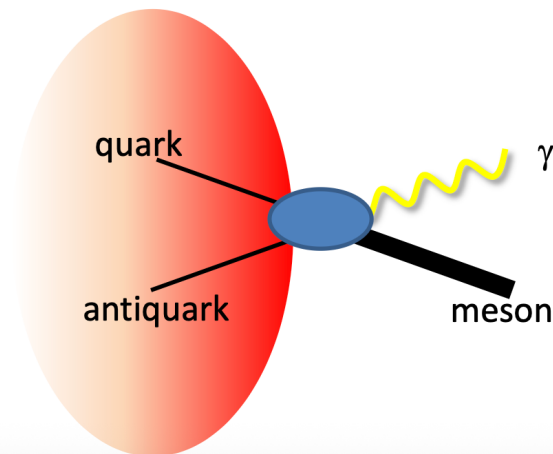
Radiative Recombination in QGP

- Non perturbative process
- No inverse process in HIC
- Non equilibrium process

⇒ Construction of dynamical model

Recombination Model

スクリーンショット

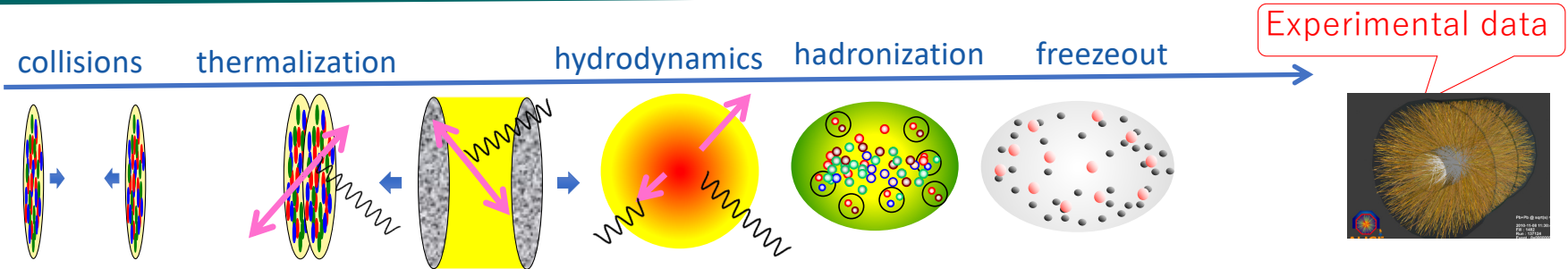


C. NUNAKA

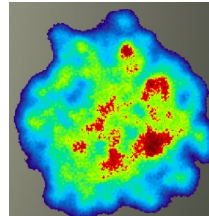


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From Fluid to Particle



Initial
conditions



Energy (entropy)
density distributions

Hydrodynamics

QCD phase
diagram
EoS: lattice QCD
Shear and bulk
viscosities

Final state interactions

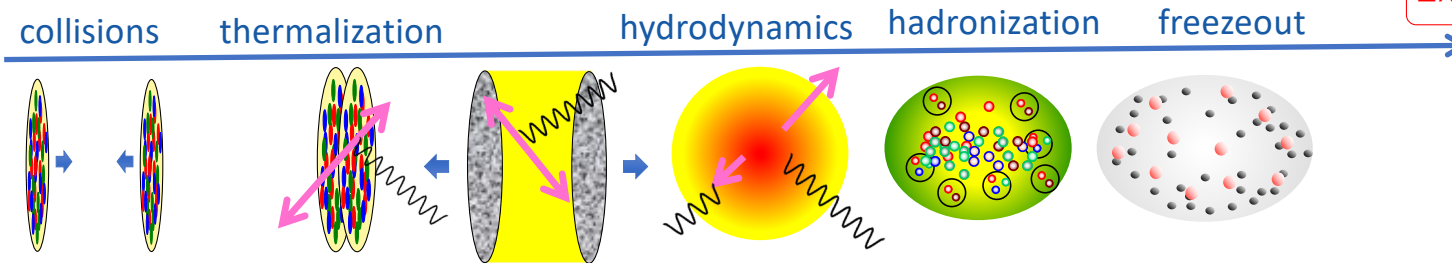
Hadron based event
generator

Hadronization

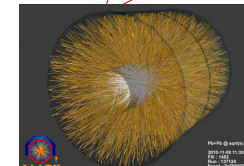
Recombination Model



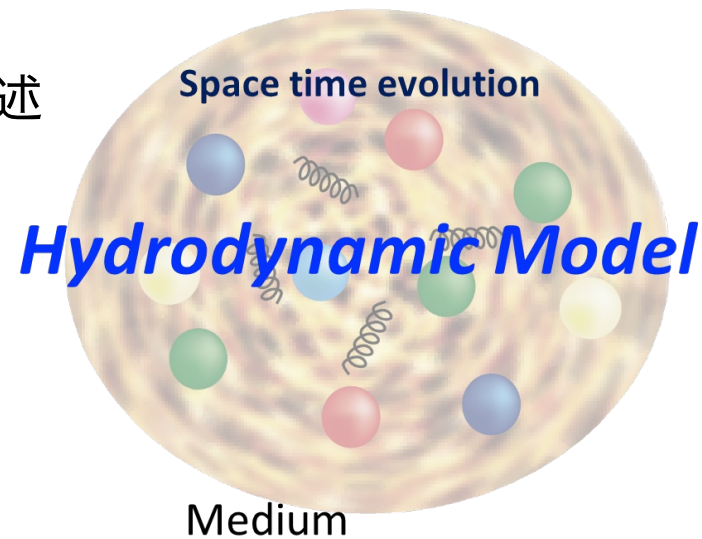
QGP物理の理解



Experimental data



ダイナミクスの記述



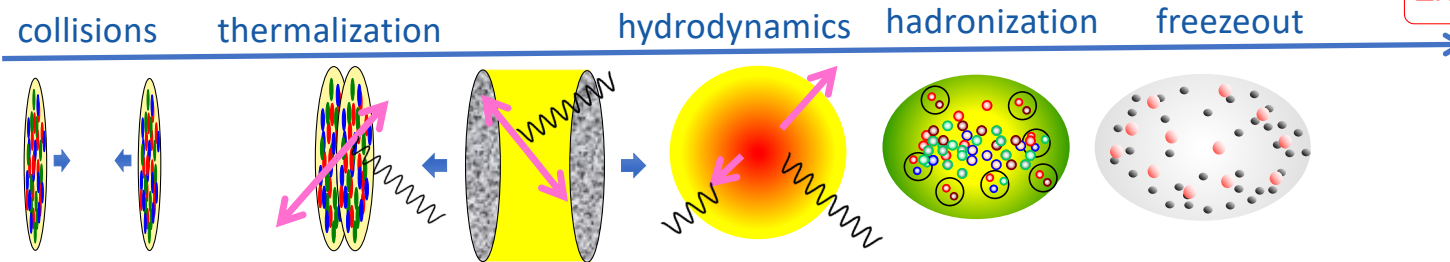
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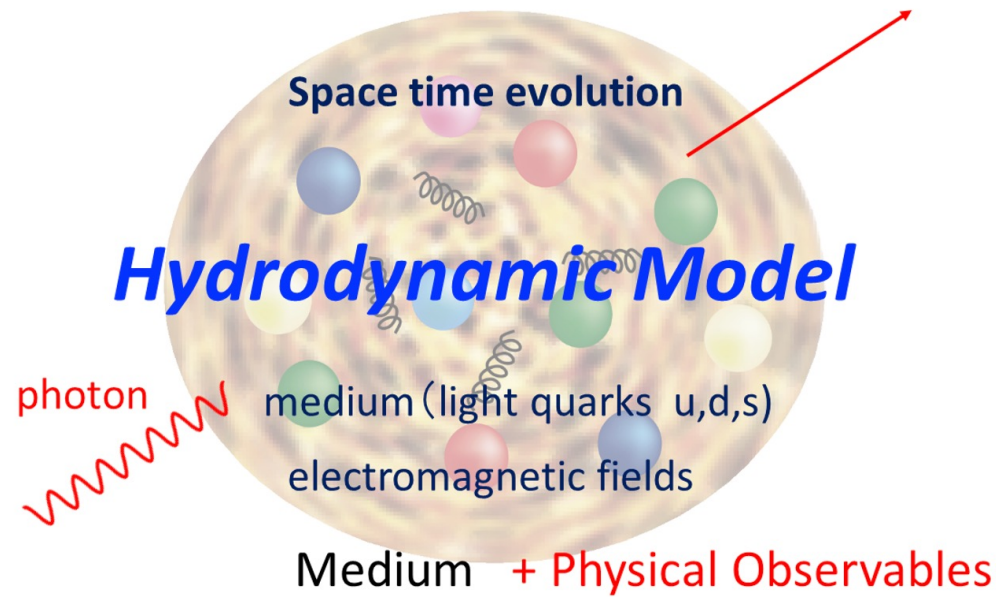
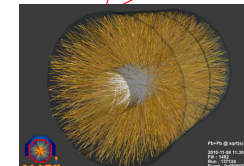
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QGP物理の理解



Experimental data

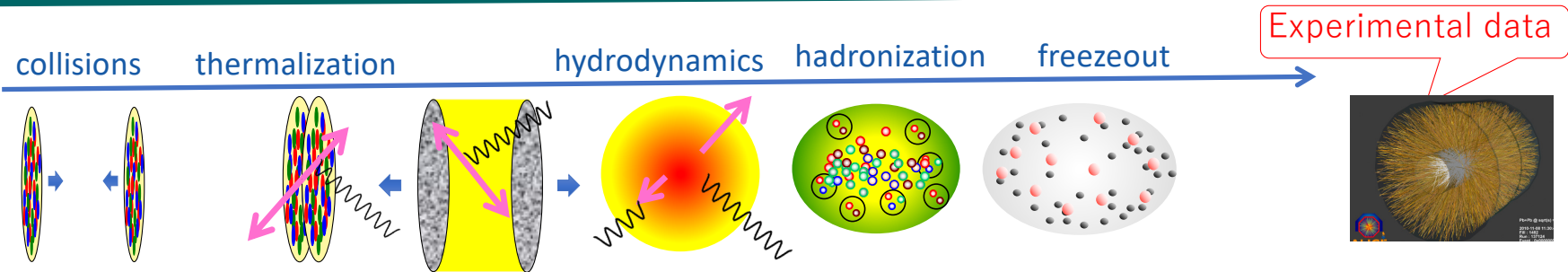


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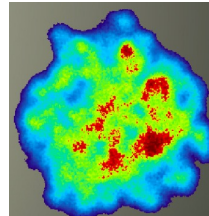


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From Fluid to Particle



Initial
conditions



Energy (entropy)
density distributions

TRENTO

Hydrodynamics

QCD phase
diagram
EoS: lattice QCD
Shear and bulk
viscosities

New
hydrodynami
cs code

$$\partial_\mu T^{\mu\nu} = 0$$

Final state interactions

Hadron based event
generator

Hadronization

Cooper-Frye Formula

Akamatsu et al, JCP256,34(2014)

Okamoto, Akamatsu, Nonaka,

EPJC76,579(2016)

Okamoto and Nonaka, EPJC77,383(2017)

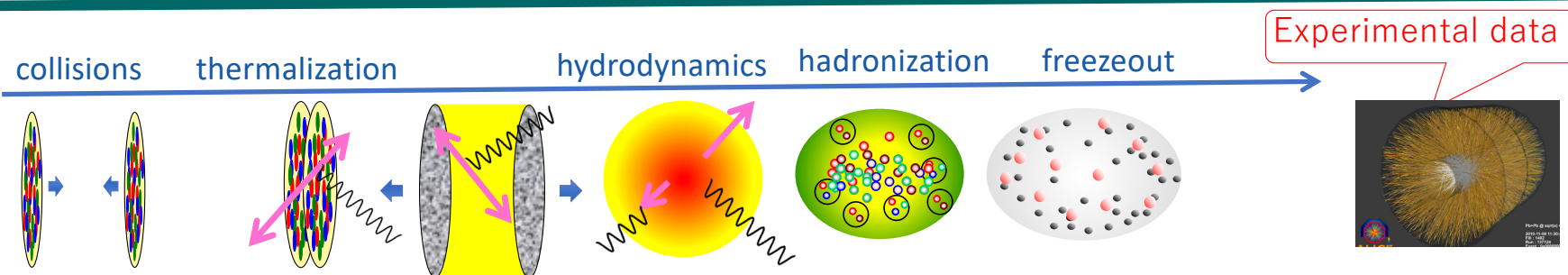


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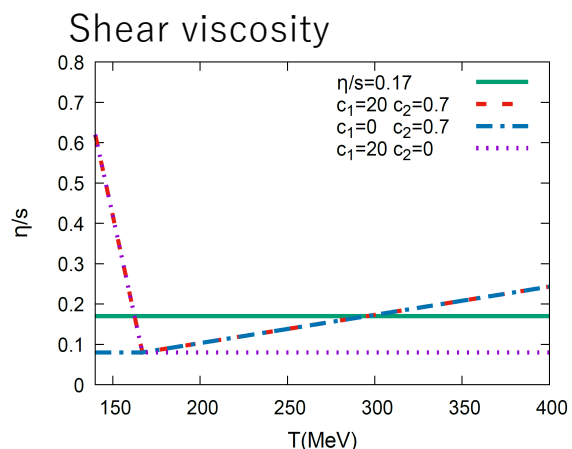
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From Fluid to Particle



Our Model

Experimental data



Bulk viscosity

$$\zeta = b\eta \left(\frac{1}{3} - c_s^2 \right)^2 \quad b = 40$$

ALICE Pb+Pb $\sqrt{s_{NN}} = 2.76$ TeV, LHC

- ✓ Rapidity distributions
- ✓ P_T distributions
- ✓ Mean P_T
- ✓ Collective flows v_2 and v_3

What physical observable is interesting?



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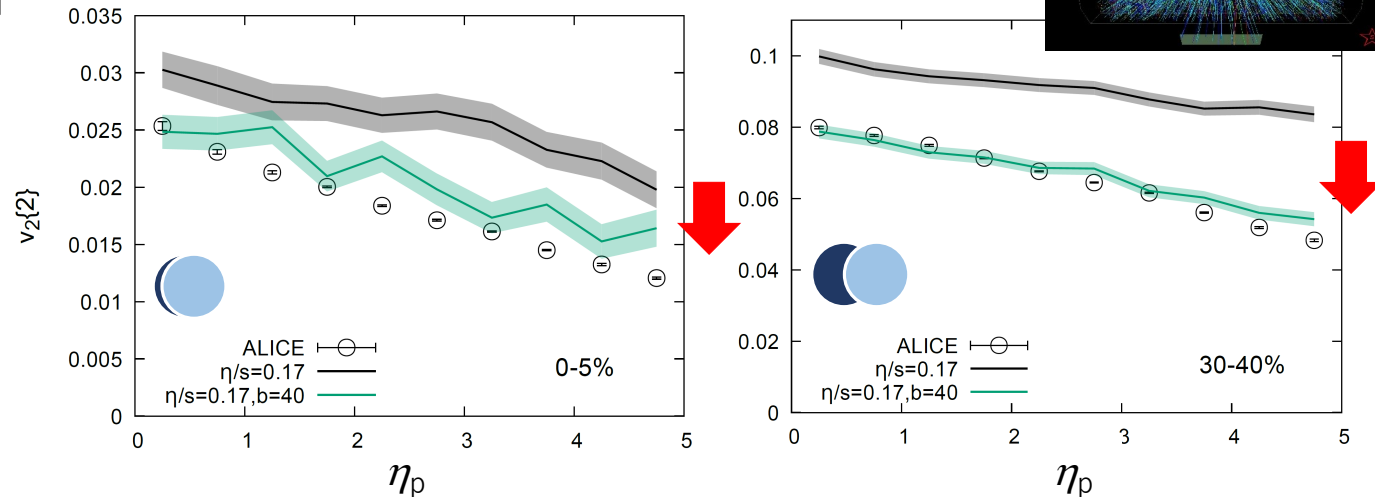


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Effect on Collective Flow



Collective flow as a function of pseudorapidity



- (3+1)-d calculation
- v_n with bulk viscosity is much closer to the ALICE data: amplitude and slope
- Effect of bulk viscosity at forward rapidity is large.

Finite bulk viscosity



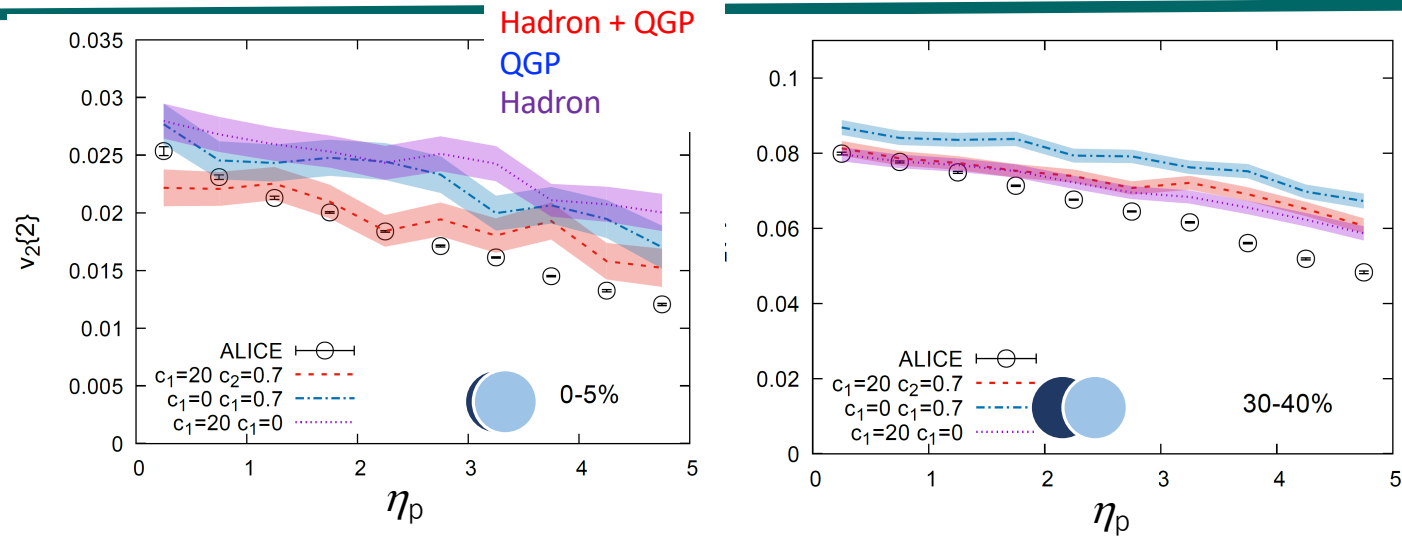
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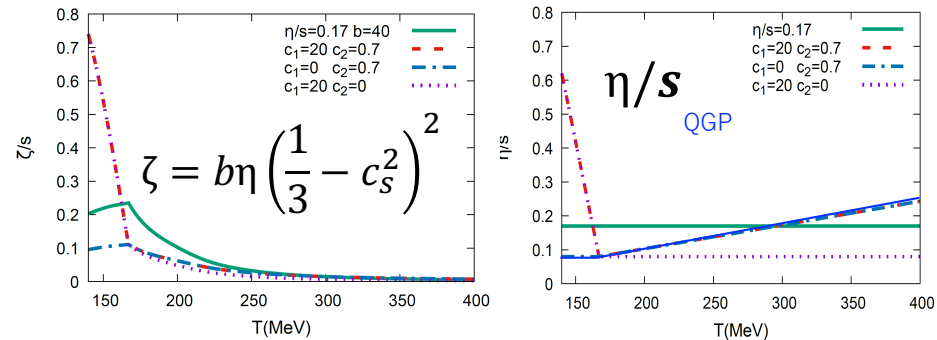
32

Temperature Dependent η/s



QGP phase: $\eta/s(T)$
Hadron phase: $\eta/s=0.08$

In both centrality classes
 $v_2(\eta_p)$ is larger.



$$T_{\text{SW}} = 150 \text{ MeV}$$

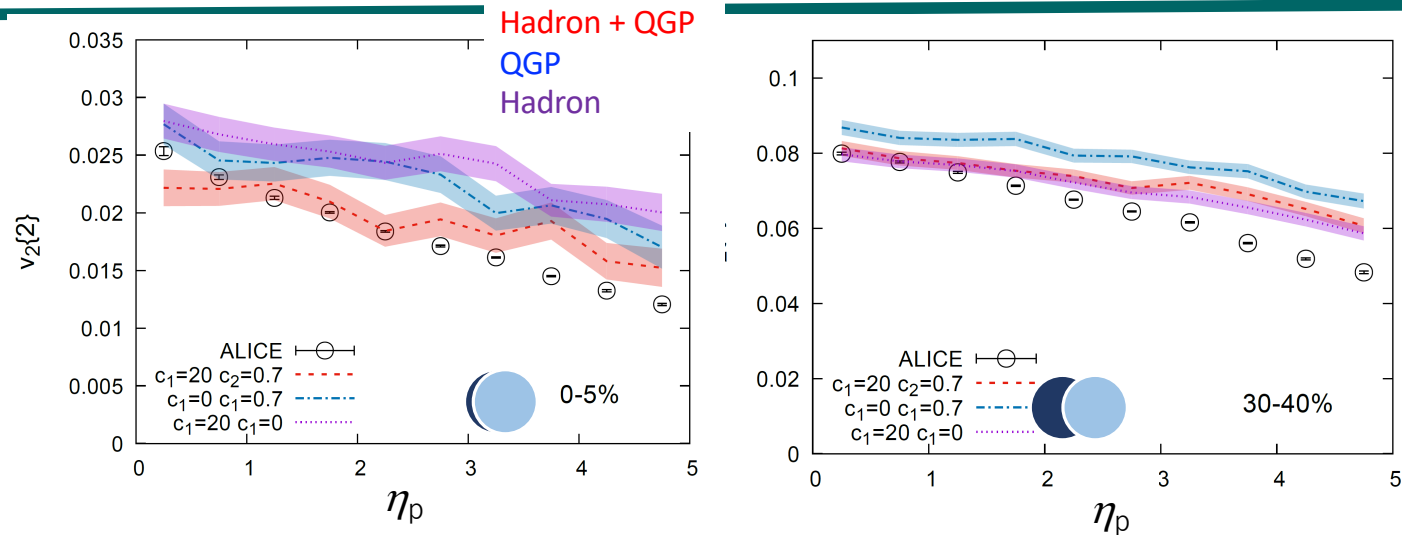


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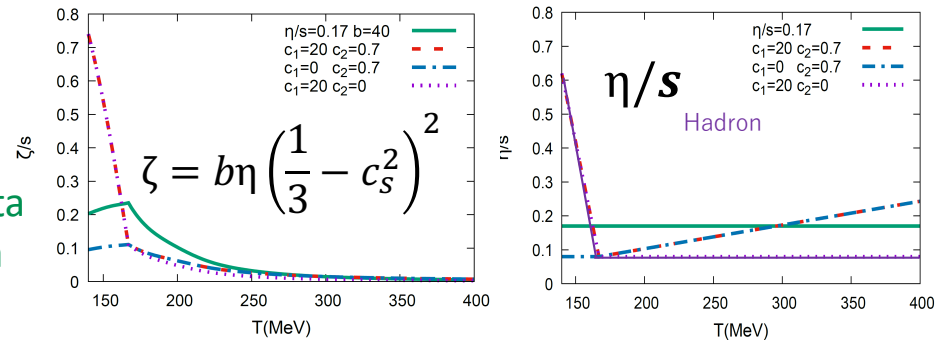
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Temperature Dependent η/s



QGP phase: $\eta/s=0.08$
 Hadron phase: $\eta/s(T)$

0-5%: larger than ALICE data
 30-40%: close to ALICE data



$$T_{\text{SW}} = 150 \text{ MeV}$$

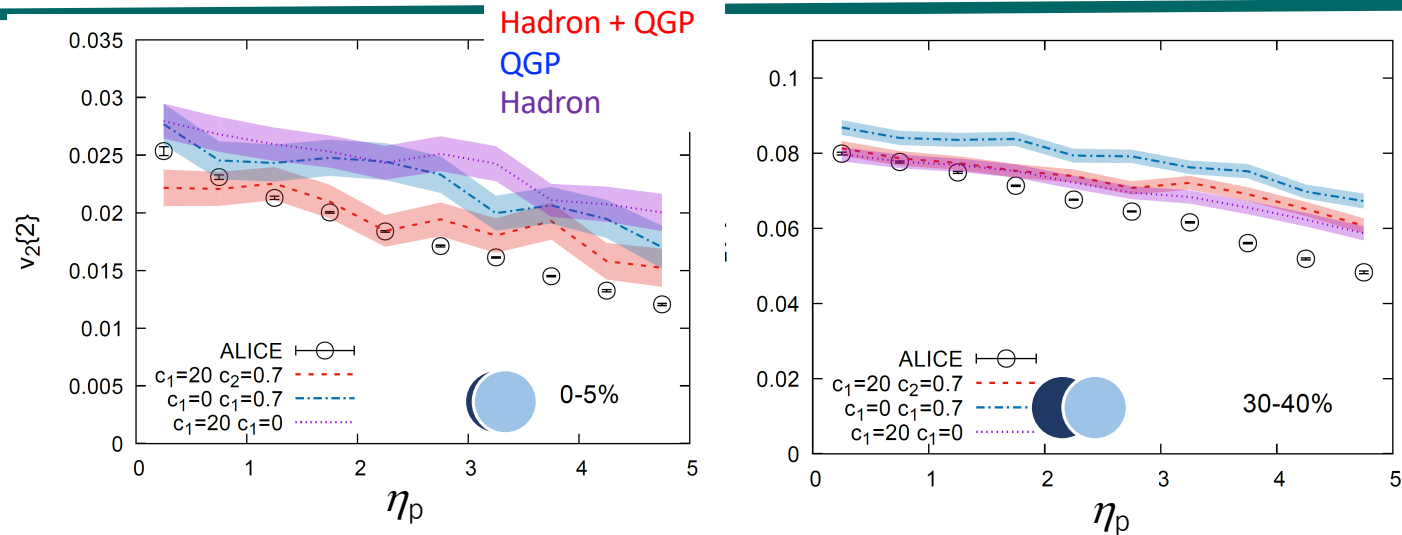


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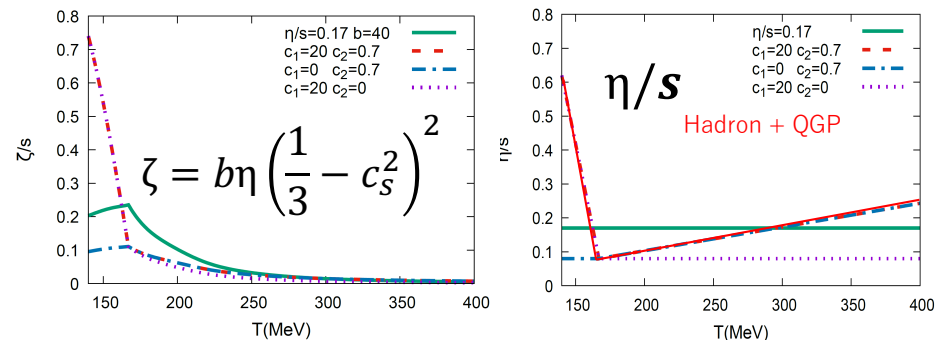
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Temperature Dependent η/s



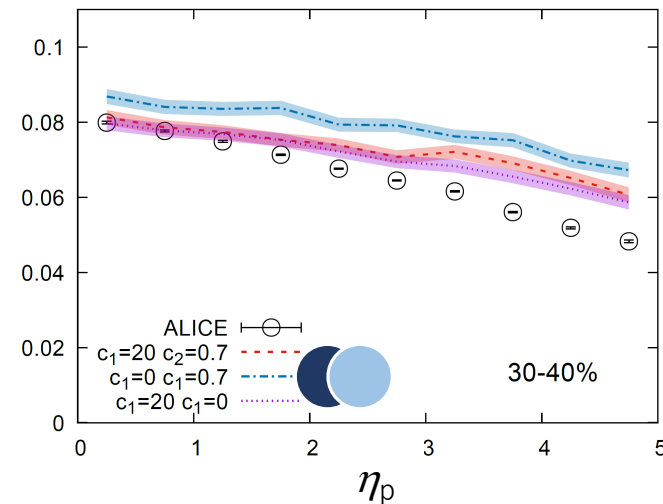
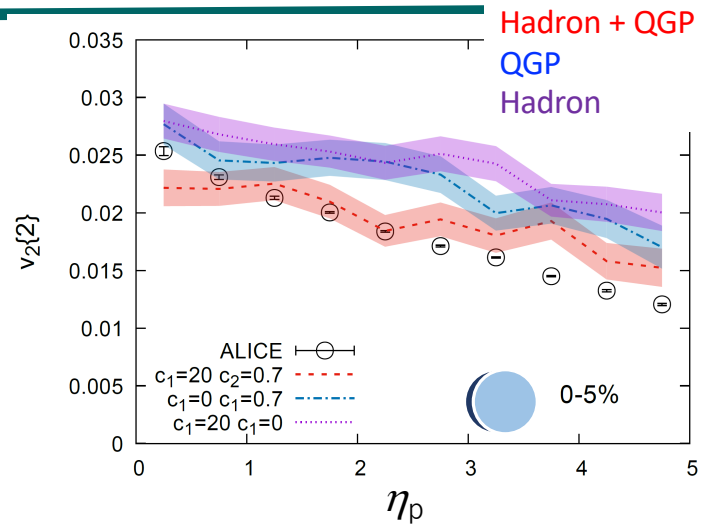
QGP phase: $\eta/s(T)$
Hadron phase: $\eta/s(T)$

In both central classes,
close to ALICE data



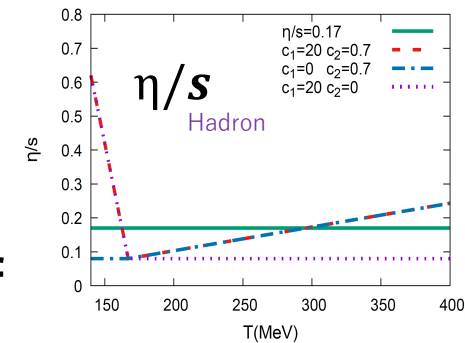
$$T_{\text{SW}} = 150 \text{ MeV}$$

Temperature Dependent η/s



- 0-5 % centrality
 η/s of QGP and hadron phases is important.
- 30-40 % centrality
 η/s of hadron phase is important.

Central dependence of $v_2(\eta_p)$
reveals temperature dependence of
viscosities.

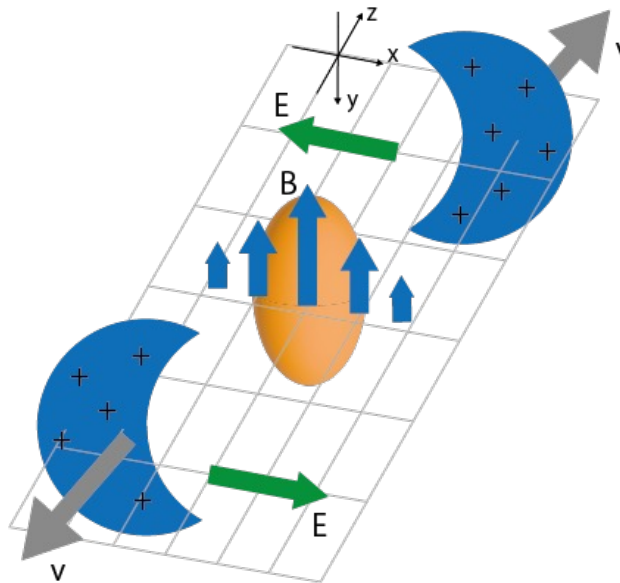


Electromagnetic Fields in Heavy Ion Collisions ?



- **Strong Electromagnetic field ?**

- Au + Au ($\sqrt{s_{NN}} = 200$ GeV) : 10^{14} T $\sim 10 m_{\pi}^2$
- Pb + Pb ($\sqrt{s_{NN}} = 2.76$ TeV) : 10^{15} T



Electromagnetic Fields and Property of QGP



• Electric Conductivity

• Dissipation of electric field

- Ampere's law : $\partial_t \vec{E} - \nabla \times \vec{B} = -\vec{j}$

$\vec{B}$: magnetic field
 $\vec{E}$: electric field

Ohm's law makes electric field dissipate

→ Dissipated energy to fluid (medium)

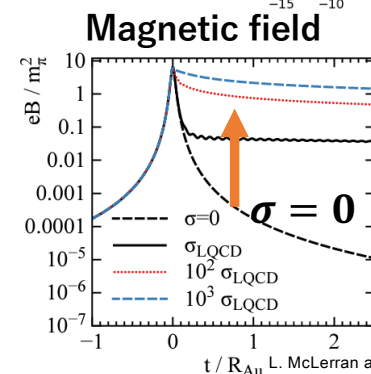
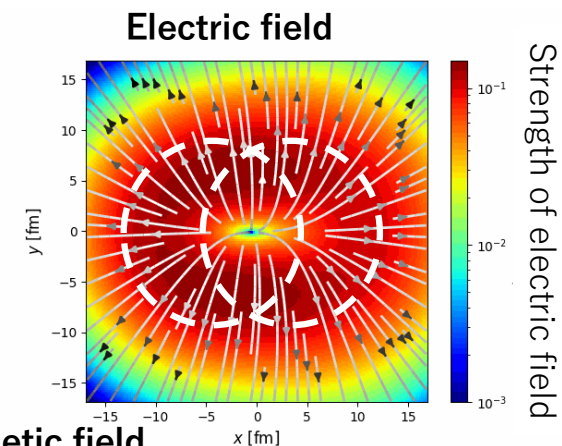
• Charge is induced by electric field

- Induced charge depends on charge conductivity

• Dissipation of magnetic field

Charge conductivity of QGP

← dissipation of electromagnetic fields and charge distribution QGP



$\sigma \neq 0$

Suppresses of dissipation

Electric field is dissipated.

L. McLerran and V. Skokov, Nucl. Phys. A 929 (2014), 184-190



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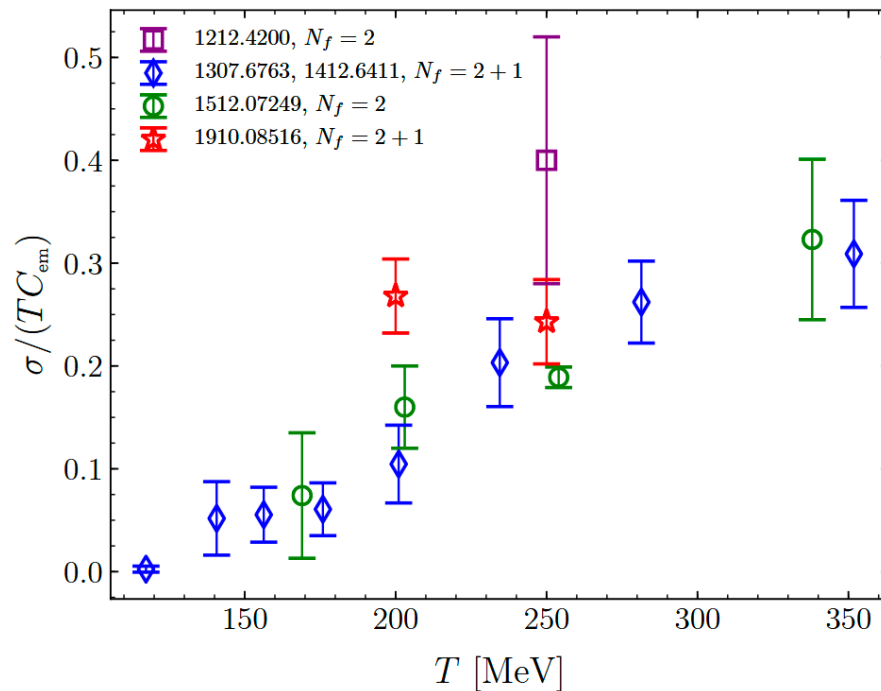


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Electric Conductivity of QCD Matter



• Lattice QCD



Aarts, Nikolaev, *EPJ.A* 57, 118 (2021); 2008.12326 [hep-lat]

Electric Conductivity on the Lattice

$$\sigma = \frac{1}{6} \frac{\partial}{\partial \omega} \left(\int d^4x e^{i\omega t} \langle [j_\mu^{\text{em}}(t, x), j_\mu^{\text{em}}(0, 0)] \rangle \right) \Big|_{\omega=0}$$

Uses linear-response theory (Kubo formula)

Low energy limit of the electromagnetic spectral function

- Does not include external magnetic field effects
- Uses approximately realistic pion mass
- General agreement among results using a variety of methods and parameters

Relativistic resistive magnetohydrodynamics



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Understanding of QGP Property



Charge conductivity of QGP from analysis of high-energy heavy-ion collisions

Physical property	Observables	Quantitative analysis
Charge conductivity	Charge dependent flow, EM probes..	Just started
Shear viscosity	Azimuthal anisotropy v_n	○
Bulk viscosity	P_T distributions	○
Diffusion coefficient	Jet energy loss	○

Charge dependent directed flow

Asymmetric collisions → i.e., Hirono, Hongo, and Hirano, PRC 90, 021903 (2014).

Symmetric collisions

Proposed EM observables

Dileptons → i.e., Akamatsu, Hamagaki, Hatsuda, and Hirano, PRC 85, 054903 (2012).
Photons → i.e., Sun and Yan, PRC 109, 034917 (2024).

Construction of relativistic resistive magnetohydrodynamics

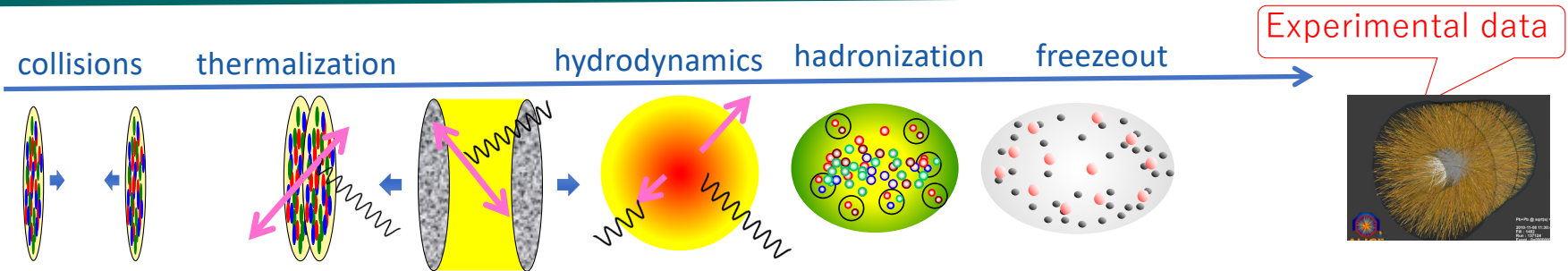


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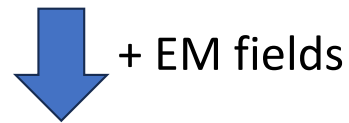
Space-Time Evolution of HIC



Initial condition

Hydrodynamics

Final state interactions



Initial condition

Magnetohydrodynamics

Final state interactions



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Smoothed Initial Condition



Tuchin, Phys.Rev.C88,024911(2013)

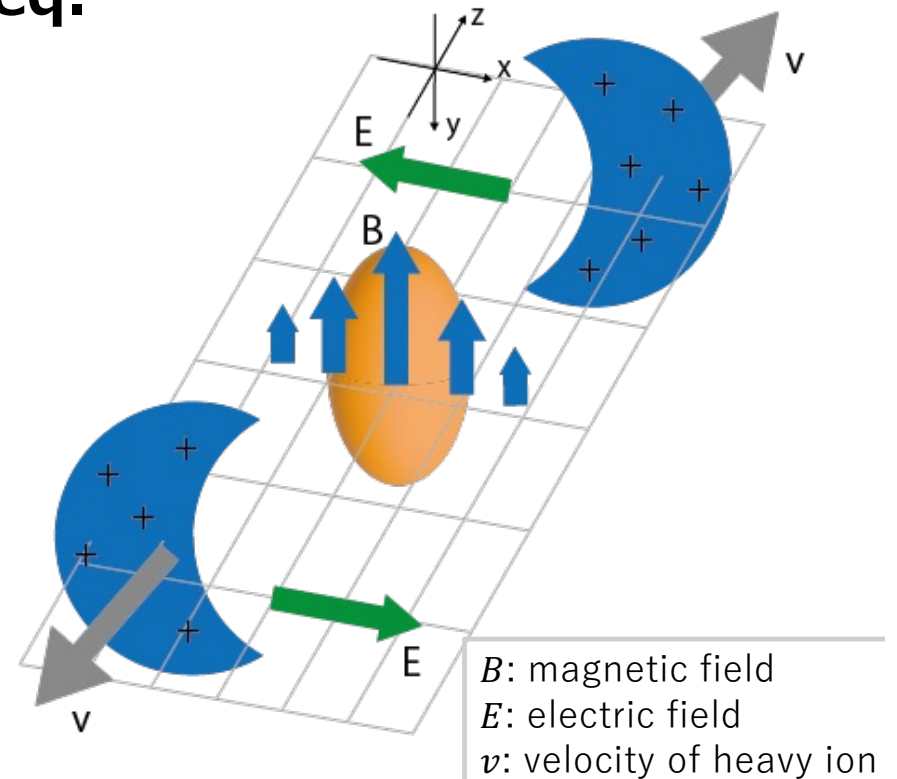
■ Asymptotic solution of Maxwell eq.

➤ Electromagnetic field made by point charge moving in the longitudinal axis

- Proton distribution in nucleus : uniform sphere
- Constant charge conductivity ($\sigma = 0.023 \text{ fm}^{-1}$)

$$\begin{aligned}\nabla \cdot \mathbf{B} &= 0, & \nabla \times \mathbf{E} &= -\frac{\partial \mathbf{B}}{\partial t}, \\ \nabla \cdot \mathbf{D} &= e\delta(z - vt)\delta(\mathbf{b}), \\ \nabla \times \mathbf{H} &= \frac{\partial \mathbf{D}}{\partial t} + \sigma \mathbf{E} + ev\hat{z}\delta(z - vt)\delta(\mathbf{b})\end{aligned}$$

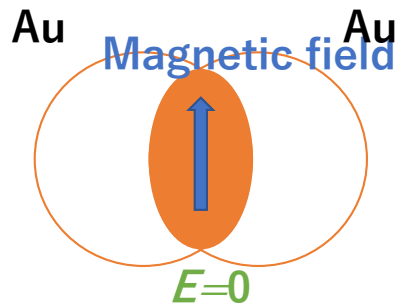
Integration of the asymptotic solutions over the charge distribution inside of nucleus



Electromagnetic Field in Symmetric and Asymmetric Systems

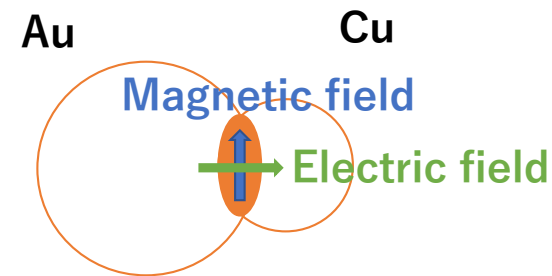


■ Au-Au collisions



- Magnetic field
 - Strong magnetic field
- Electric field
 - No electric field

■ Cu-Au collisions



- Magnetic field
 - Strong magnetic field
- Electric field
 - $E \neq 0$ due to the asymmetry of the charge distribution

Hirono, Hongo, Hirano

Initial Condition : QGP Medium



Bozek, et al, Phys. Rev. C 81, 054902(2010)

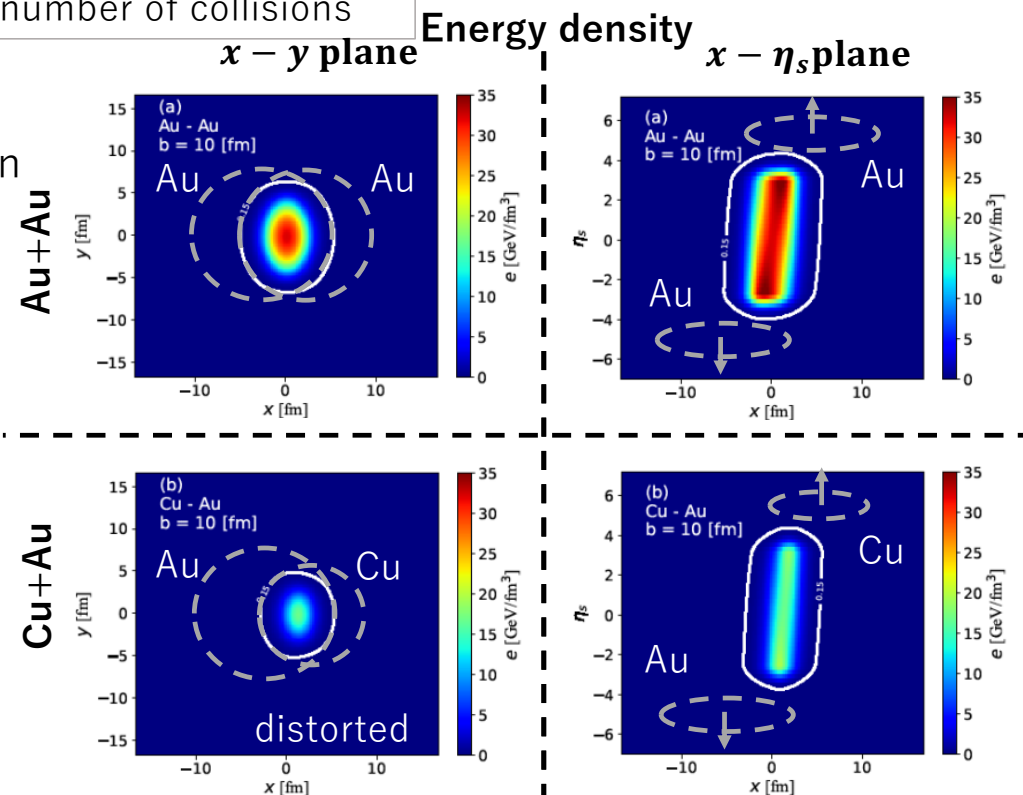
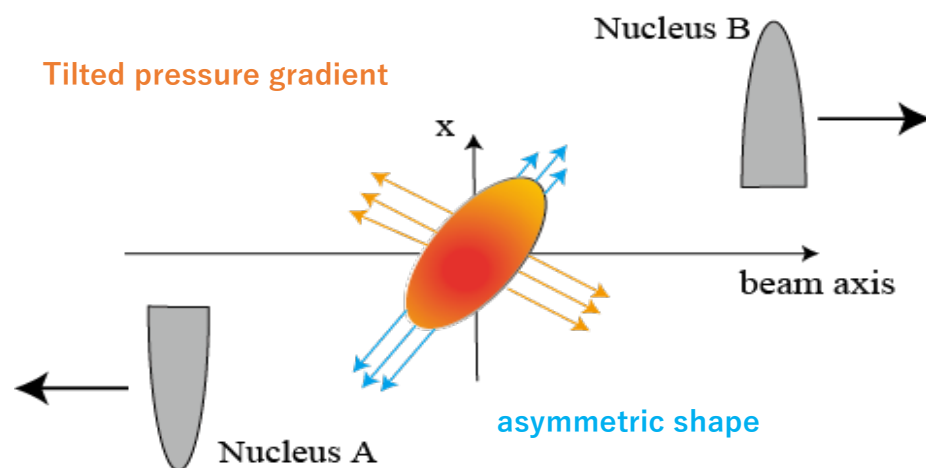
Freezeout hypersurface

■ Tilted Glauber model

n_p : number of participants
 n_c : number of collisions

- Energy density is scaled by n_p and n_c
- Tilted distribution in the longitudinal direction

For directed flow v_1



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Initial Condition : Electromagnetic Fields ($\eta_s = 0$)



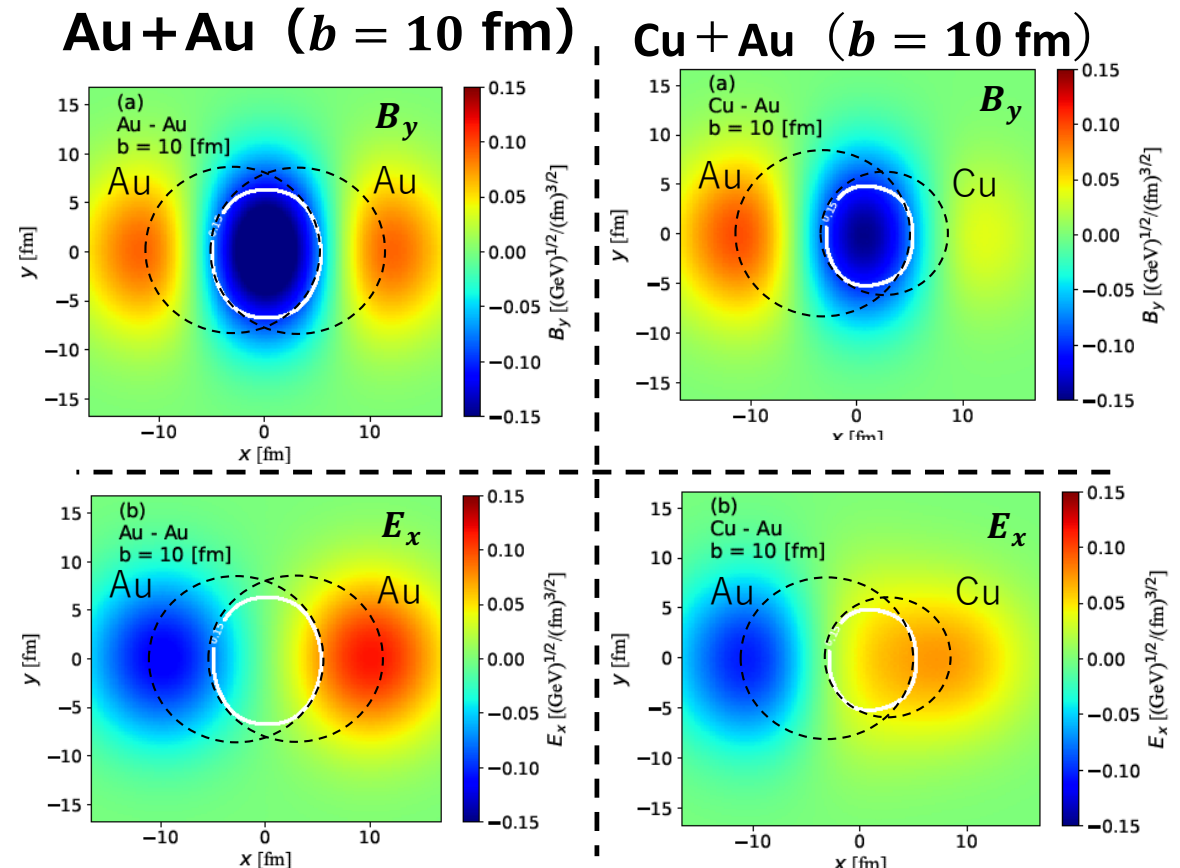
Tuchin, Phys.Rev.C88,024911(2013)

■ Au+Au

- Strong magnetic fields in QGP
- Electric field ~ 0 in QGP

■ Cu+Au

- Strong magnetic field in QGP
- Finite electric field in QGP



Freezeout hypersurface

Magnetohydrodynamics



• Magnetohydrodynamics

- Conservation law

$$\nabla_\mu T^{\mu\nu} = 0$$



$$\nabla_\mu T_m^{\mu\nu} = -J_\mu F^{\mu\nu}$$

Energy momentum tensor

$$T^{\mu\nu} = T_m^{\mu\nu} + T_f^{\mu\nu}$$

matter

$$T_m^{\mu\nu} = (e + p)u^\mu u^\nu + pg^{\mu\nu}$$

Maxwell eq.

$$\nabla_\mu F^{\mu\nu} = -J^\nu$$

EM field

$$T_f^{\mu\nu} = F^{\mu\lambda} F_\lambda^\nu - \frac{1}{4} g^{\mu\nu} F^{\lambda\kappa} F_{\lambda\kappa}$$

$$\nabla_\mu T_f^{\mu\nu} = J_\mu F^{\mu\nu}$$

- Ohm's law *Blackman and Field, PRL71,3481(1993)*

$$J^\mu = \sigma F^{\mu\nu} u_\nu + qu^\mu$$

$$q = -J^\mu u_\mu$$

Electrical conductivity



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Magnetohydrodynamics



• Magnetohydrodynamics

- Conservation law

$$\nabla_\mu T^{\mu\nu} = 0$$



$$\nabla_\mu T_m^{\mu\nu} = -J_\mu F^{\mu\nu}$$

Energy momentum tensor

$$T^{\mu\nu} = T_m^{\mu\nu} + T_f^{\mu\nu}$$

matter

$$T_m^{\mu\nu} = (e + p)u^\mu u^\nu + pg^{\mu\nu}$$

+viscosity, Hattori..

Maxwell eq.

$$\nabla_\mu F^{\mu\nu} = -J^\nu$$

EM field

$$T_f^{\mu\nu} = F^{\mu\lambda} F_\lambda^\nu - \frac{1}{4} g^{\mu\nu} F^{\lambda\kappa} F_{\lambda\kappa},$$

$$\nabla_\mu T_f^{\mu\nu} = J_\mu F^{\mu\nu}$$

- Ohm's law *Blackman and Field, PRL71,3481(1993)*

$$J^\mu = \sigma F^{\mu\nu} u_\nu + qu^\mu$$

$$q = -J^\mu u_\mu$$

Charge diffusion, *Dash et al., PRD107,056003(2023)*

Electrical conductivity



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Magnetohydrodynamics



• RRMHD equation

➤ Conservation law + Maxwell eq. + Ohm's law

$$\partial_\mu T^{\mu\nu} = F^{\nu\lambda} J_\lambda$$

$$J^\mu = J_c^\mu + qu^\mu$$

e : energy density

p : pressure

$$p_{em} = (E^2 + B^2)/2$$

$$\varepsilon = (e + p)\gamma^2 - p + p_{em}$$

$$m^i = (e + p)\gamma^2 v^i + \epsilon^{ijk} B_j E_k$$

$$\Pi^{ij} = (e + p)\gamma^2 v^i v^j + (p + p_{em})g^{ij} - E^i E^j - B^i B^j$$

Energy Conservation

$$\partial_t \varepsilon + \nabla \cdot m = 0$$

Momentum conservation

$$\partial_t m^i + \nabla \cdot \Pi^i = 0$$

Faraday's law

$$\partial_t \vec{B} + \nabla \times \vec{E} = 0$$

Ohm's law

$$\vec{J} = q\vec{v} + \sigma\gamma[\vec{E} + \vec{v} \times \vec{B} - (\vec{v} \cdot \vec{E})\vec{v}]$$

Ampere's law

$$=: \vec{J}_c$$

$$\partial_t \vec{E} - \nabla \times \vec{B} = q\vec{v}$$

$$\partial_t \vec{E} = \vec{J}_c$$

operator
splitting

- Integration with quasi-analytic solutions

$$\vec{E}_\perp = -\vec{v} \times \vec{B} + (E_\perp^0 + \vec{v} \times \vec{B}) \exp(-\sigma\gamma t)$$

$$\vec{E}_\parallel = E_\parallel^0 \exp(-\sigma t/\gamma)$$

Benoît, Miyoshi, C. N., and Takahashi, ArXiv: 2502.04611

Nakamura, Miyoshi, C. N. and Takahashi, Phys. Rev. C 107, (2023) 014901

Nakamura, Miyoshi, C. N. and Takahashi, Eur.Phys.J.C 83 (2023) 3, 229.

Nakamura, Miyoshi, C. N. and Takahashi, Phys. Rev. C 107 (2023) 3, 034912



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RRMHD Equation in Milne Coordinates



New

- Milne coordinates
 - Expanding systems in the longitudinal direction (τ, x, y, η_s)
 - Strong expansion in the longitudinal direction is effectively included.
 - Number of grid of fluid is saved.

$$\tau = \sqrt{t^2 - z^2}$$

$$\eta_s = \frac{1}{2} \ln \frac{t+z}{t-z}$$

RRMHD Equation

$$\partial_\tau(\tau U) + \partial_i(\tau F^i) = \tau S$$

$$U = \begin{pmatrix} D \\ m_j \\ \varepsilon \\ B^j \\ E^j \\ q \end{pmatrix}, F^i = \begin{pmatrix} Dv^i \\ \Pi^{ji} \\ m^i \\ \varepsilon^{jik} E_k \\ \varepsilon^{jik} B_k \\ J^i \end{pmatrix}, S = \begin{pmatrix} 0 \\ \frac{1}{2} T^{ik} \partial_j g_{ik} \\ -\frac{1}{2} T^{ik} \partial_0 g_{ik} \\ 0 \\ J_c^i \\ 0 \end{pmatrix}$$

The first RRMHD code in Milne coordinates



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Validation of the Code

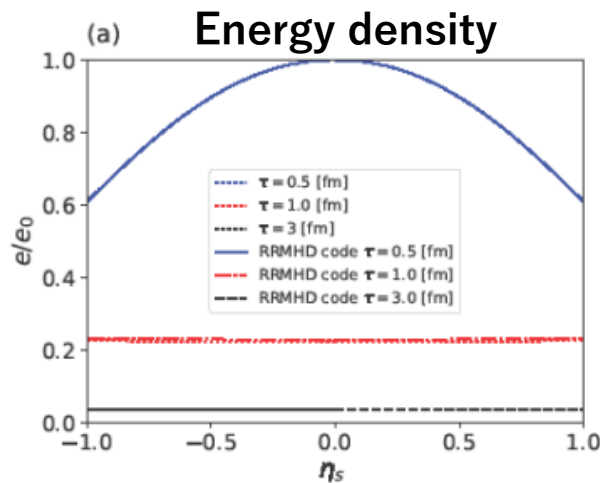


- RRMHD in the Milne coordinates

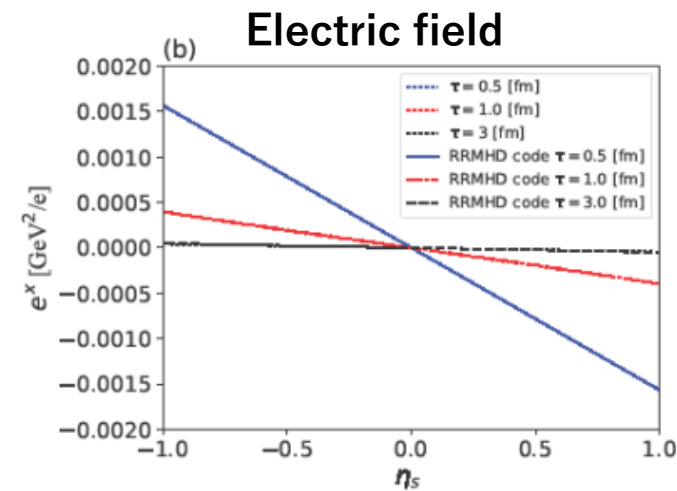
Nakamura, Miyoshi, CN and Takahashi, Eur.Phys.J.C 83 (2023) 3, 229.

New Test Problem

- (1+1) dimensional expansion system $u^\mu = (\cosh Y, 0, 0, \sinh Y)$
 - Comparison between quasi-analytical solution and RRMHD simulation



Solid line : RRMHD code
Dashed line: quasi-analytical solution



➔ Application to Heavy Ion Collisions



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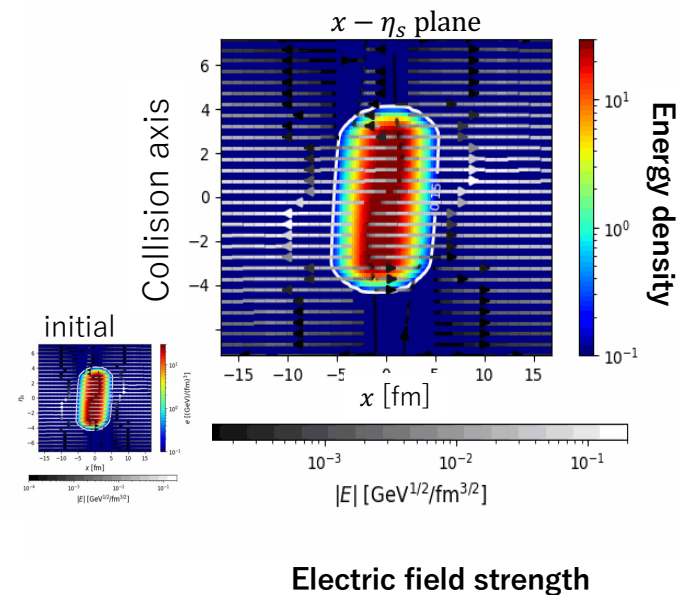
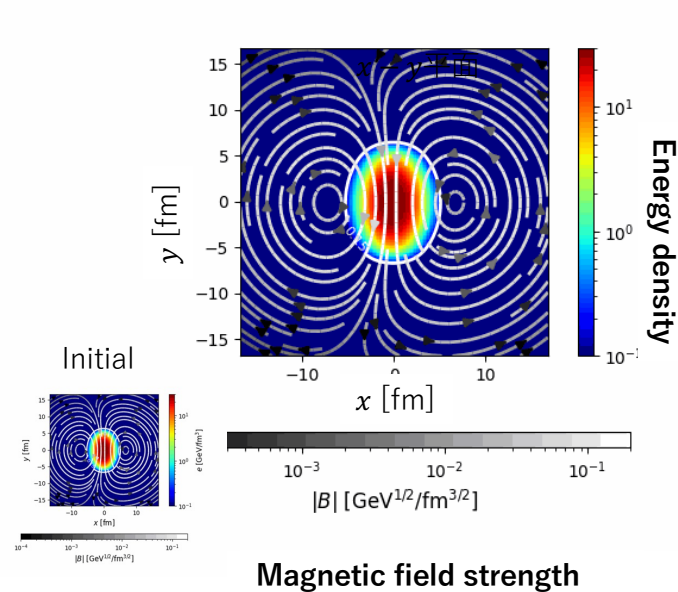
Space-time Evolution



Nakamura, Miyoshi, CN and Takahashi, PRC 107, no.1, 014901 (2023)

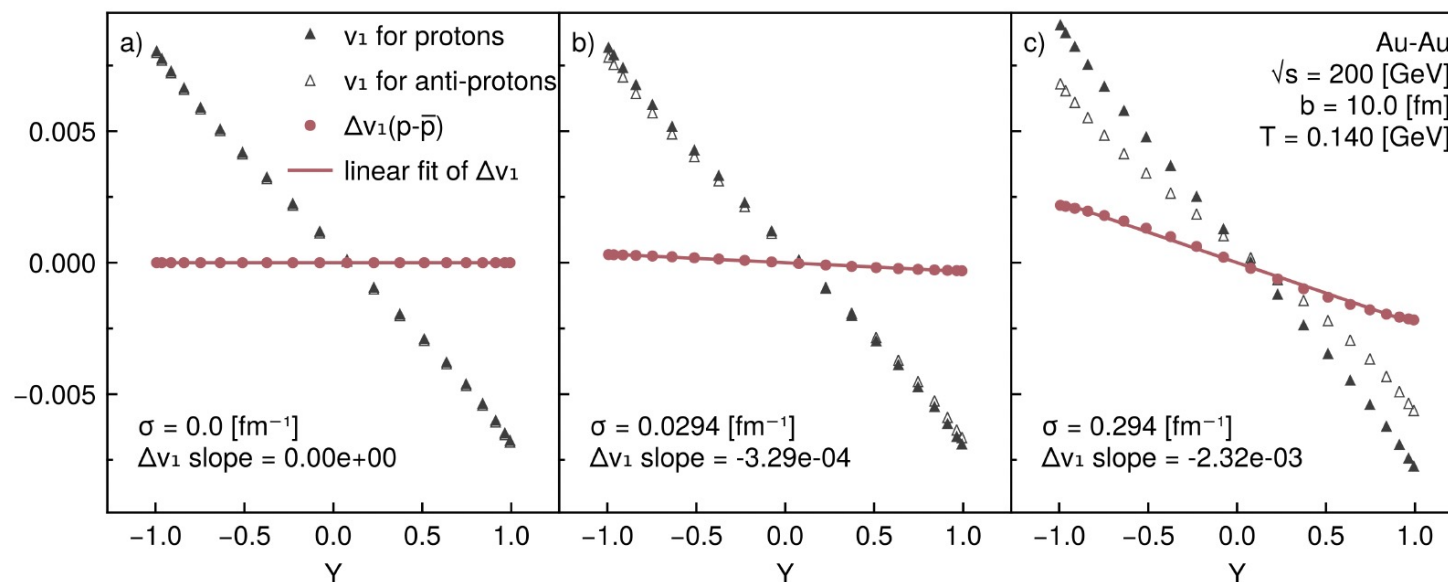
Au+Au collision system

First calculation in HIC with RRMHD code



Analysis of Heavy Ion Collisions

Charge Dependent Directed Flow



Benoit, Miyoshi, C. N., and Takahashi, ArXiv: 2502.04611

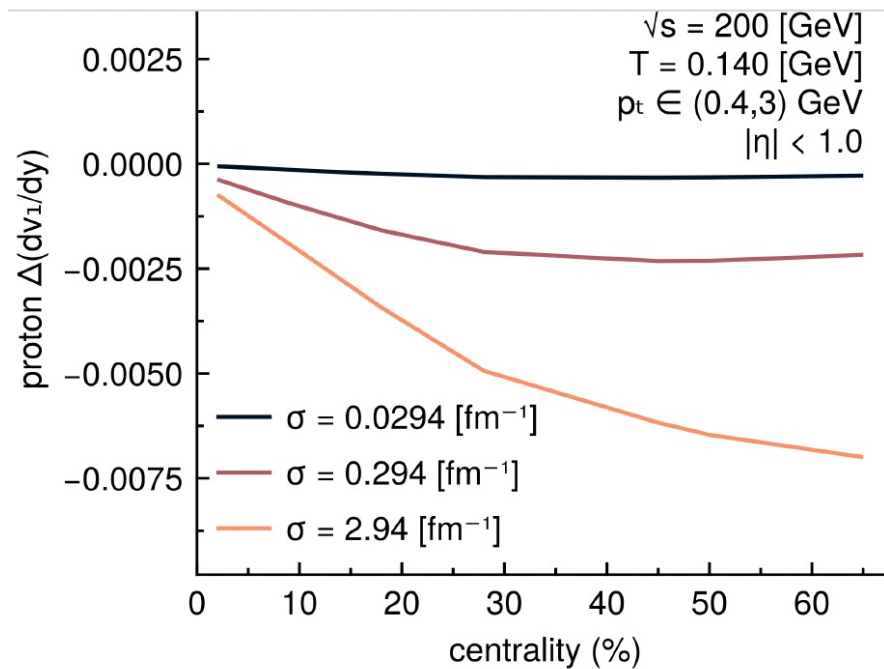


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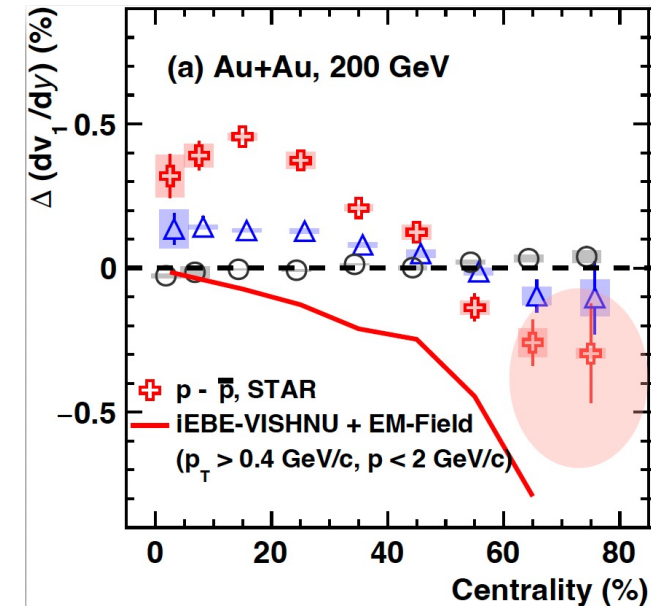


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Charge Dependent Flow



Benoit, Miyoshi, C. N., and Takahashi, ArXiv: 2502.04611



Caveat:

No baryon current

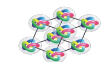
Initial condition

EoS

Final state interactions

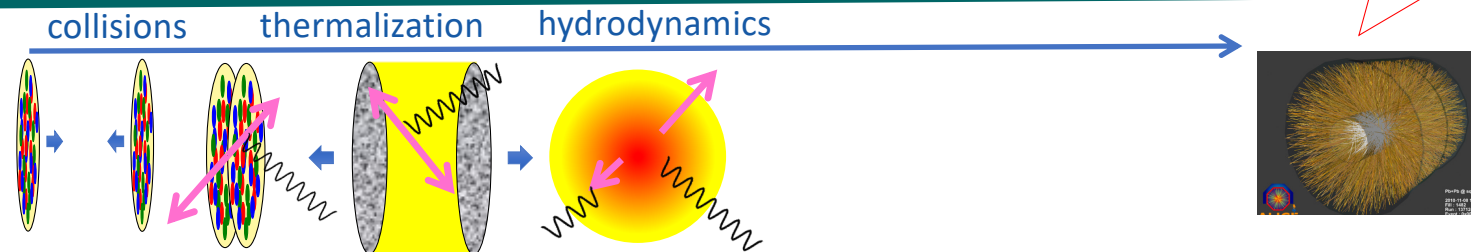


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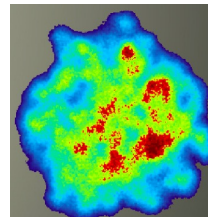


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Initial Condition



Initial
conditions



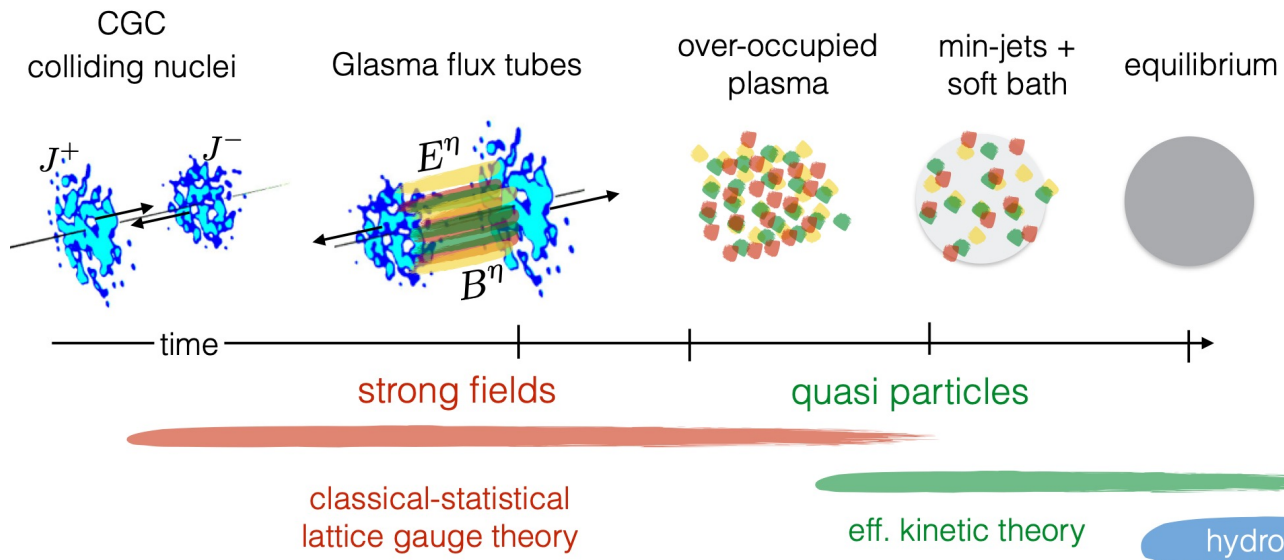
Energy (entropy)
density distributions

最も難しい問題の一つ

Hydrodynamics

QCD phase diagram
EoS: lattice QCD
Shear and bulk viscosities

General picture at weak-coupling



Consistent theoretical description requires combination of different weak-coupling methods

3

Schlichting [Initial Stages 2016]



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Initial Condition (Thermalization)

Abdi



- Transport model : approximation solution to Boltzmann Transport Equation.

$$\frac{\partial}{\partial t} f(t, \vec{r}, \vec{p}) + \frac{\vec{p}}{m} \nabla_{\vec{r}} f(t, \vec{r}, \vec{p}) - \nabla_{\vec{r}} U(\vec{r}) \nabla_{\vec{p}} f(t, \vec{r}, \vec{p}) = \left(\frac{\partial f}{\partial t} \right)_{coll}$$

Time evolution

Diffusion

External field

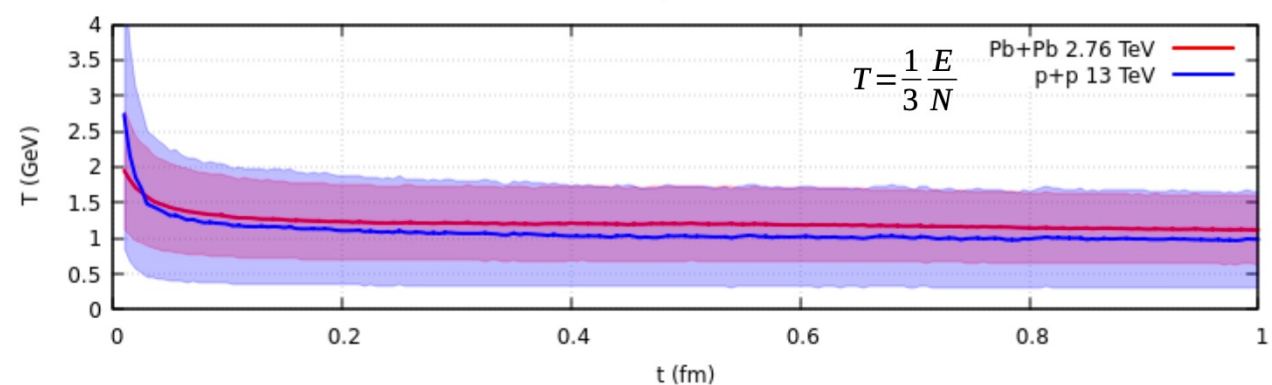
Interparticle collision

Parton Cascade Model

- Applicable to far-from-equilibrium quark-gluon plasma.
- Bottom up approach to quark-gluon plasma thermal and chemical equilibration.
- Based on Hadronic transport model SMASH^[1] reconfigured for quark-gluon plasma case.

Goal : Create a realistic initial condition generator model for hydrodynamic simulation with the minimum number of free parameters.

[1] J. Weil, V. Steinberg, J. Staudenmaier, L. G. Pang, D. Oliinych, Goldschmidt, B. Bäuchle, J. Auvinen, M. Attems, and H. Petersen



- Temperature is constantly decreasing due to the expanding medium
- Saturation behaviour can be seen well before 0.2 fm for both collision systems



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Initial Conditions : Quark Production

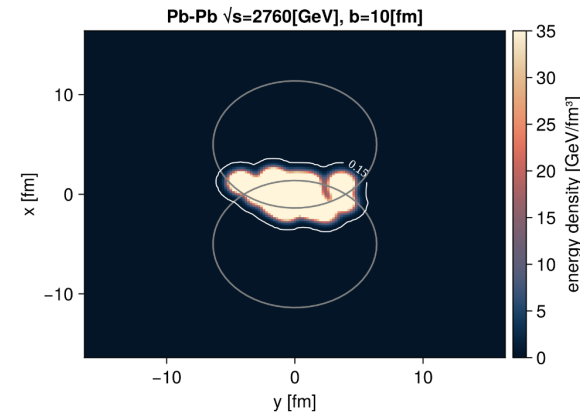
Benoit, Taya



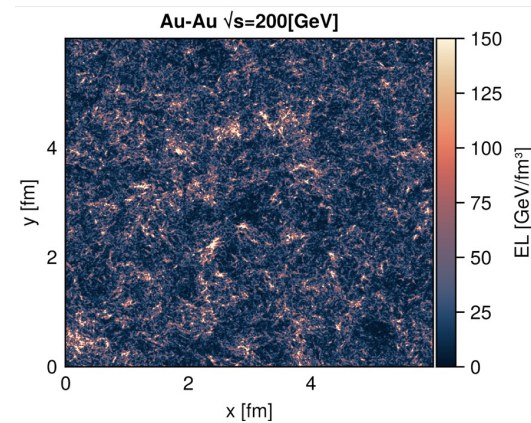
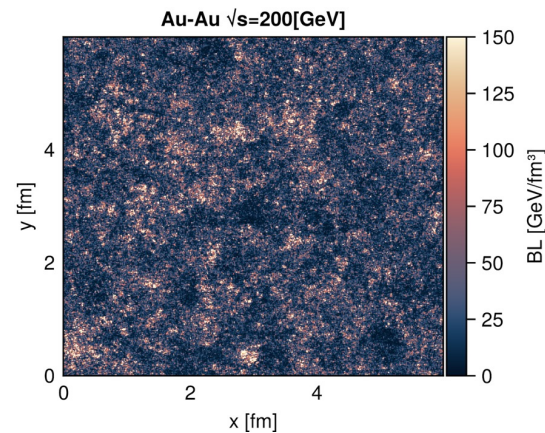
- Smoothed initial conditions
- Fluctuating initial conditions
 - Smaller B field, Larger E field
- Glasma initial condition

Gluon場中の粒子

Particle in Cell



Benoit

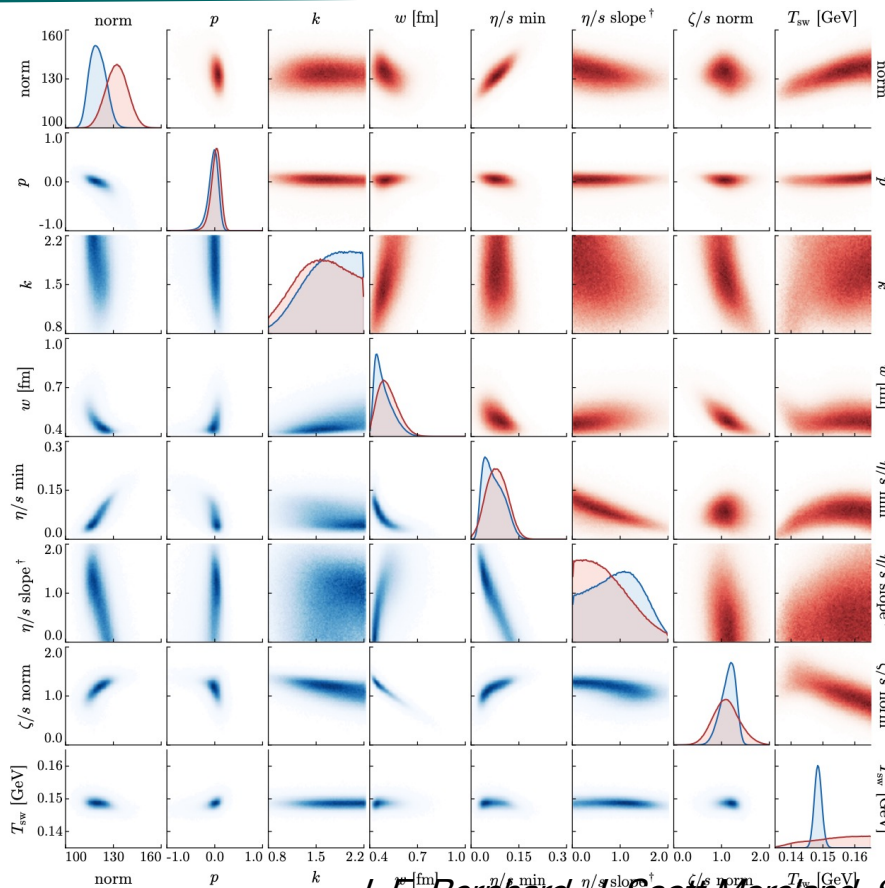


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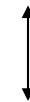


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ベイジアン解析



流体模型中の
パラメータ



複数の実験結果



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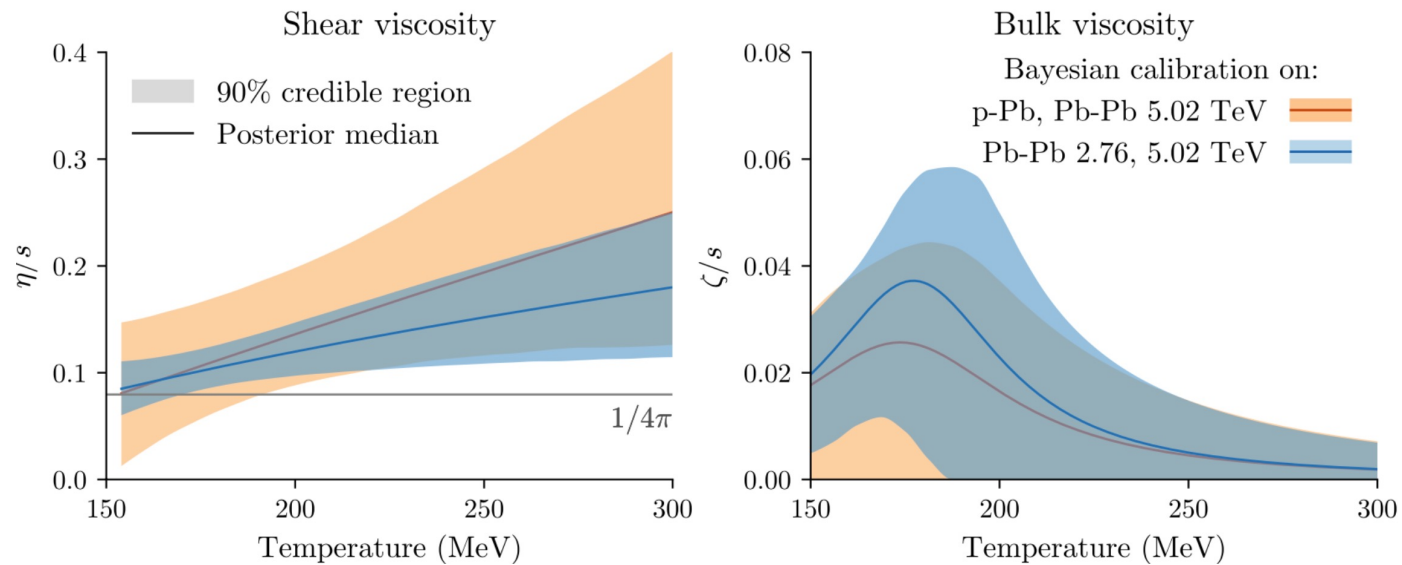
J. E. Bernhard, J. Scott Moreland, S. A. Bass, J. Liu, U. Heinz, PRC 94.024907(2016)



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粘性の温度依存性



あくまでも模型内での結果

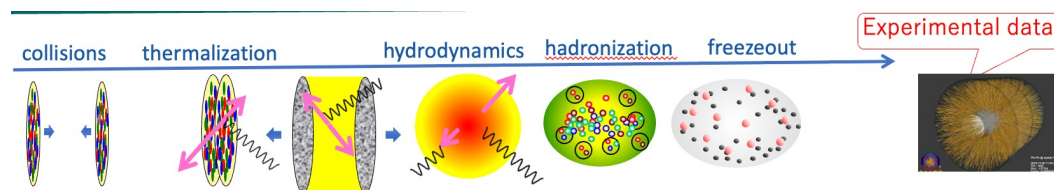
クォーク・グルーオン プラズマの物理



- 世界規模の実験が稼働中：理論と実験の検証

- 周辺物理との関連

- 理論的概念、ツール：共通点
- 基礎理論
 - 相転移、相図
 - 平衡、非平衡物理
 - 熱力学量
- ダイナミクスの記述
 - 相対論的流体模型
 - 分子動力学
 - Particle in Cell
- データ解析、大規模計算
 - ベイジアン解析



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